

Phys 661 HW 2 Solutions

Problem 1 – John Daniel

1. The 6 isotopes of Thorium that were used:

→ Thorium-217: Decay time of 0.241 ms

Energy of α -decay: 9261 keV

Note: When a number of different α decays

were possible with accompanying probabilities, the

mean energy of the α -decay was used.

→ Thorium-221: Decay time of 1.73 ms

Energy of α -decay: 8251 keV

→ Thorium-223: Decay time of 0.6 s

Energy of α -decay: 7298 keV

→ Thorium-226: Decay time of 30.57 minutes

Energy of α -decay: 6300 keV

→ Thorium-228: Decay time of 1.9125 years

Energy of α -decay: 5368.976 keV

→ Thorium-230: Decay time of 7.54×10^4 years

Energy of α -decay: 4657.378 keV

→ The data for the time was then converted to units of years

& a plot of $\ln\left(\frac{\tau_{1/2}}{\text{years}}\right)$ vs $\frac{1}{\sqrt{E_\alpha}}$, where E_α was in units of eV was created.

→ The result was a fairly linear graph with a slope of approximately 281847.13 ± 15695.3 .

→ If the Thorium-217 data point is ignored, the slope, or the proportionality constant between $\ln\left(\frac{\tau_{1/2}}{\text{years}}\right)$ and $\frac{1}{\sqrt{E_\alpha}}$, is 302867 ± 74686 .

→ If we then compare this with the WKB Approximation, we see it is fairly consistent.

→ The WKB Approximation predicts the relationship is:

$$\ln \frac{T_{1/2}}{1 \text{ year}} \approx 2\pi(Z-2)\alpha_e \sqrt{\frac{2M_\alpha c^2}{E}}$$

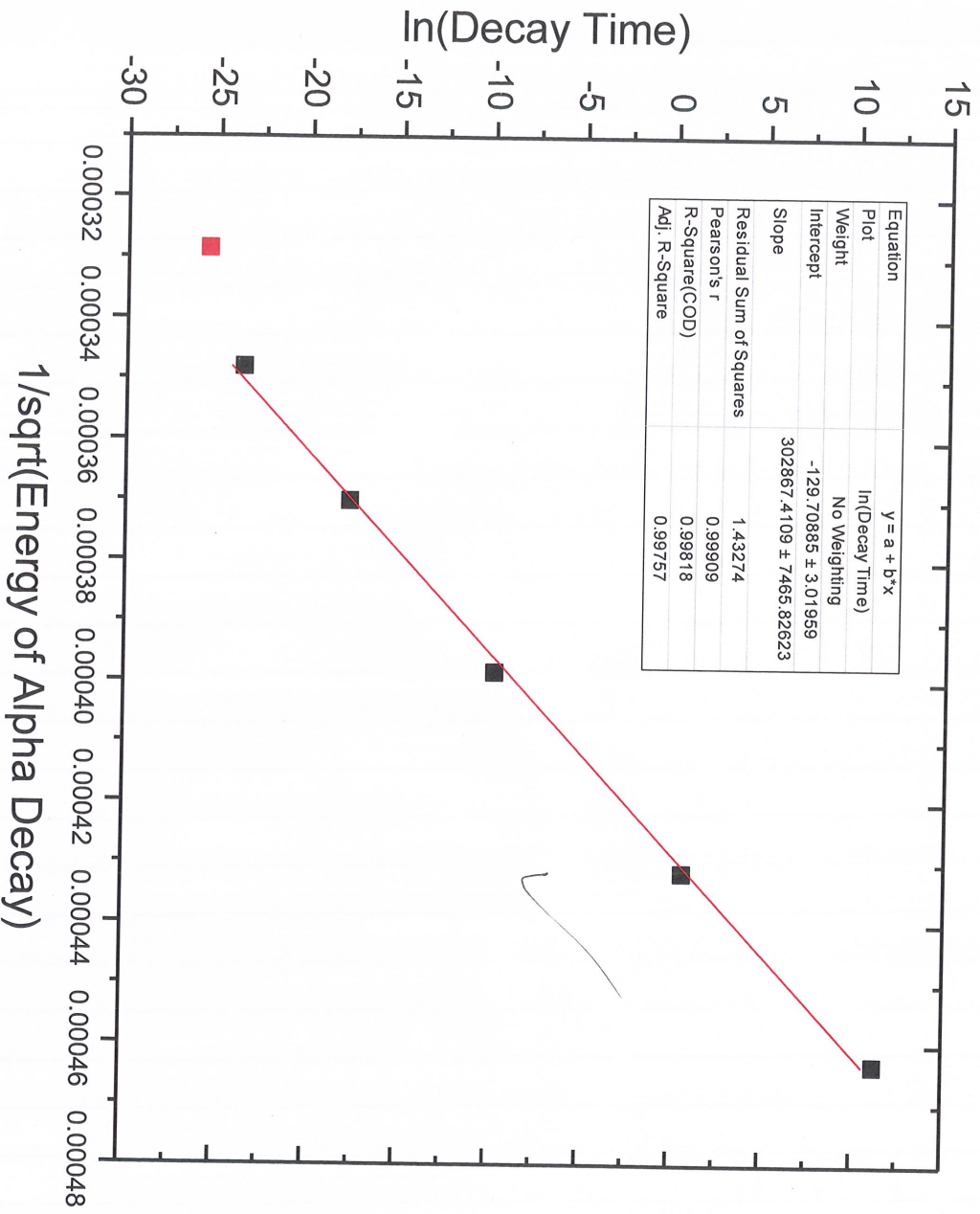
where (for Thorium) $Z=90$, $\alpha_e = \frac{1}{137}$ (the fine structure constant) & $M_\alpha = 6.64424 \times 10^{-27} \text{ kg}$ (the mass of an α particle)

→ When this is evaluated (& ensuring to divide the $M_\alpha c^2$ term by 1.6×10^{19} to get the $M_\alpha c^2$ term in terms of eV), we find the slope \ln is predicted to be ≈ 348932.0728

∴ The WKB approximation gives

$$\ln \left(\frac{T_{1/2}}{1 \text{ year}} \right) \approx 348932.0728 \cdot \frac{1}{\sqrt{E_\alpha}}, \text{ where } E_\alpha \text{ is in terms of eV}$$

⇒ We see that this agrees broadly with what is observed (as is, it is the same order of magnitude), but there is a difference of around 4000.



1/sqrt(Energy of Alpha Decay)

Problem 2 – Sean Myers

Problem 2

For Parity $\Pi |l l_z\rangle = (-)^l |l l_z\rangle$.

If you have a matrix element

$$\langle \alpha | V | \beta \rangle$$

then it equals $\langle \alpha | \Pi \Pi V \Pi \Pi | \beta \rangle$

$$\langle \alpha | V | \beta \rangle = \epsilon_\alpha \epsilon_\beta \epsilon_V \langle \alpha | V | \beta \rangle$$

if $\epsilon_\alpha \epsilon_\beta \epsilon_V = -1$ the element

Also $\Pi |r\rangle = |-r\rangle$; so $\Pi |z\rangle = \cancel{1} |z\rangle$ must be zero.

from table: parity gives

$$\#1 \quad \langle 1,0 | z | 2,0 \rangle \Rightarrow (-)^1 (-)^2 (-) = + \quad (\text{no info gained}) \checkmark$$

$\uparrow \quad \uparrow$
 $\Pi \Pi \quad \Pi \Pi$

$$\#2 \quad \langle 1,0 | z y^2 | 2,1 \rangle = (-)^1 (-)^2$$

$\uparrow \quad \uparrow$
 $\Pi \Pi \quad \Pi \Pi$

$$(-)^1 (-)^2 \langle 1,0 | \underbrace{\Pi z y^2 \Pi}_{-1} | 2,1 \rangle \Rightarrow + \quad (\text{no info}) \checkmark$$

$$\#3 \quad \langle 4,0 | p_z \times | 0,0 \rangle = (-)^4 (-)^0 (-)(-). \Rightarrow + \quad (\text{no info}) \checkmark$$

$$\#4 \quad \langle 3,0 | p_x^2 | 0,0 \rangle = (-)^3 (-)^0 (+) \Rightarrow - \quad (\text{zero}) \checkmark$$

$$\#5 \quad \langle 2,1 | x p_x | 1,0 \rangle = (-)^2 (-)^1 (-)(-) \Rightarrow - \quad (\text{zero}) \checkmark$$

Now for L^2 and L_z we use

$\langle l, m_1 | T_q^k | l_2, m_2 \rangle$ vanishes if
the following ~~is~~ is not satisfied:

for L^2 : $|l_2 - k| \leq l_1 \leq l_2 + k$

L_z : $m_1 = q + m_2$

First I determine T_q^k for each operator

#1 $z = T_0^1 \quad k=1 \quad q=0$

$\langle 1, 0 | T_0^1 | 2, 0 \rangle \quad m_1 = m_2 = 0$
 $l_1 = 1; l_2 = 2$

L^2 : $|2-1| \leq 1 \leq 2+1$
 $1 \leq 1 \leq 3$ true (no info)

L_z : $0 = 0 + 0$ true (no info)

#2 $zy^2 \approx T_0^1 y^2$

$y \approx T_{+1}^1 + T_{-1}^1$

$y^2 \approx T_{+1}^1 T_{+1}^1 + T_{-1}^1 T_{-1}^1 + T_{+1}^1 T_{-1}^1 + T_{-1}^1 T_{+1}^1$

~~$y^2 \approx T_{+1}^1 T_{+1}^1 + T_{-1}^1 T_{-1}^1 + T_{+1}^1 T_{-1}^1 + T_{-1}^1 T_{+1}^1$~~

$zy^2 = T_0^1 T_{+1}^1 T_{+1}^1 + T_0^1 T_{-1}^1 T_{-1}^1 + T_0^1 T_{+1}^1 T_{-1}^1 + T_0^1 T_{-1}^1 T_{+1}^1$
(a) (b) (c) (d)

checking (a) $m_1 = 0$ $m_2 = 1$ $g = 2$

$l_1 = 1$ $l_2 = 2$ $k = 3$

$$T_0^1 T_{+1}^1 T_{+1}^1 \rightarrow L^2: |2-3| \leq 1 \leq 2+3$$

1 \leq 1 \leq 5 true (no info) ✓

$$L_2: 0 = 3 \text{ false (zero)}$$

since L^2 can't give zero I won't check (b), (c), (d). However, L_2 gives zero for (a) so I will see they all give zero.

(b) $g = -2$

$$0 = -2 + 1 \neq 0 \text{ (zero)}$$

(c) $g = 0$:

$$0 = 0 + 1 \neq 0 \text{ (zero)}$$

(d) $g = 0$:

→ zero

$$\therefore L_2 \rightarrow \text{(zero)} \quad \checkmark$$

#3 $P_2 \times \propto T_0^1 (T_{+1}^1 - T_{-1}^1)$; $m_1 = 0$ $m_2 = 0$

$$\propto T_0^1 T_{+1}^1 - T_0^1 T_{-1}^1 \quad l_1 = 4 \quad l_2 = 0$$

$k=2$
 $g=1$

$k=2$
 $g=-1$

$\times L^2: |0-2| \leq 4 \leq 2$ ✓

L^2 : same as other one ✓

$L_2: 0 = 1 + 0$ false

$L_2: 0 = -1 + 0$ false

for #3 L^2, L_z give zero

$$\#4 \quad p_x^2 \propto \left(\begin{array}{cccc} T_{+1}^{\dagger} T_{+1}^{\dagger} & + T_{-1}^{\dagger} T_{-1}^{\dagger} & + T_{+1}^{\dagger} T_{-1}^{\dagger} & + T_{-1}^{\dagger} T_{+1}^{\dagger} \\ \downarrow (a) & (b) & (c) & (d) \end{array} \right)$$

$m_1 = 0 \quad m_2 = 0$
 $l_1 = 3 \quad l_2 = 0$

(a), (b), (c), (d) have same $k=2$

$$L^2: |0-2| \leq 3 \leq 0+2 \quad \times \text{ false}$$

so $L^2 \rightarrow (\text{zero})$

check (a) for $L_z: \ell = 2$

$$0 = 2+0 \neq 0 \quad (\text{no info})$$

$L_z \rightarrow (\text{no info})$ from (c) & (d)

$$\#5 \quad x p_y \propto (T_{+1}^{\dagger} - T_{-1}^{\dagger})(T_{+1}^{\dagger} + T_{-1}^{\dagger})$$

$$\propto \left(\begin{array}{cccc} T_{+1}^{\dagger} T_{+1}^{\dagger} & + T_{+1}^{\dagger} T_{-1}^{\dagger} & - T_{-1}^{\dagger} T_{+1}^{\dagger} & - T_{-1}^{\dagger} T_{-1}^{\dagger} \\ a & b & c & d \end{array} \right)$$

$$m_1 = 0 \quad m_2 = 0$$

$$l_1 = 2 \quad l_2 = 0$$

(a), (b), (c), (d) have same $k=2$

$$L^2: |1-2| \leq 2 \leq 1+3 \quad \text{true}$$

$L^2 \rightarrow \text{no info}$

$$L_z: \ell = m_1 - m_2 = 1$$

(a), (b), (c), (d) don't

have $\ell = 1 \quad \therefore L_z \rightarrow (\text{zero})$

In summary

element	Π	L^2	L_z
$\langle 1, 0 z 2, 0 \rangle$?	?	?
$\langle 1, 0 zy^2 2, 1 \rangle$?	?	0
$\langle 4, 0 p_z x 0, 0 \rangle$?	0	0
$\langle 3, 0 p_x^2 0, 0 \rangle$	0	0 ?	? 0
$\langle 2, 1 x p_x 1, 0 \rangle$	0	?	0

Problem 3 – Dmitry Kondratyev

Problem 3

To use degenerate perturbation theory we need to know the matrix elements $\langle 2l'm' | V | 2lm \rangle$. As we learned in Problem 2, operator \hat{z} adds angular momentum $|LM\rangle = |10\rangle$ so with the help of selection rules we can eliminate most of matrix elements:

$$\langle 2l'm' | V | 2lm \rangle = e|\vec{E}| \begin{pmatrix} 0 & 0 & M & 0 \\ 0 & 0 & 0 & 0 \\ M^* & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \quad \text{Where } M = \langle 210 | z | 200 \rangle.$$

|0> |1> |10> |1-1>

$$|200\rangle = \frac{1}{\sqrt{32\pi} a^{3/2}} \left[2 - \frac{r}{a} \right] e^{-r/2a}, \quad |210\rangle = \frac{1}{\sqrt{32\pi} a^{3/2}} \frac{r}{a} e^{-r/2a} \cos\theta$$

$$M = \frac{1}{32\pi a^3} \left(\int_0^\pi \cos\theta \sin\theta d\theta \right) \left(\int_0^{2\pi} d\phi \right) \left(\int_0^\infty \left[e^{-r/2a} \frac{r}{a} \left[2 - \frac{r}{a} \right] r^3 dr \right) = -3a = M^*$$

So, in the first order we have eigenstates: $|21\pm 1\rangle$, $\frac{1}{\sqrt{2}}(|200\rangle \pm |210\rangle)$
with energies: $E_2^{(0)}$, $E_2^{(0)} \mp 3ae|\vec{E}|$

Since the electric field is in \hat{z} direction, we know that polarization will be $\vec{P} = -e\vec{z} = -e\hat{z}z$. We already know that $\langle 21\pm 1 | z | 21\pm 1 \rangle = 0$, so in the eigenstates $|21\pm 1\rangle$ the polarization is 0. (equal in this case)

$$\langle \vec{P} \rangle = -e\hat{z} \cdot \frac{1}{2} (\langle 200 | \pm \langle 210 |) z (| 200 \rangle \pm | 210 \rangle) = -e\hat{z} \cdot \frac{1}{2} \left(\underbrace{\langle 200 | z | 200 \rangle}_{0} + \underbrace{\langle 210 | z | 210 \rangle}_{0} \pm \underbrace{\langle 210 | z | 200 \rangle}_{\neq 0} \pm \underbrace{\langle 200 | z | 210 \rangle}_{\neq 0} \right)$$

$$= \mp e\hat{z} \cdot \langle 200 | z | 210 \rangle = \pm e\hat{z} \cdot 3a = \pm 3ae\hat{z} \quad \text{for states } \frac{1}{\sqrt{2}}(|200\rangle \pm |210\rangle)$$

Polarization doesn't depend on $|\vec{E}|$, so it will be finite even when $E \rightarrow 0$

In the ground state polarization was proportional to $|\vec{E}|$ and was small for $E \rightarrow 0$