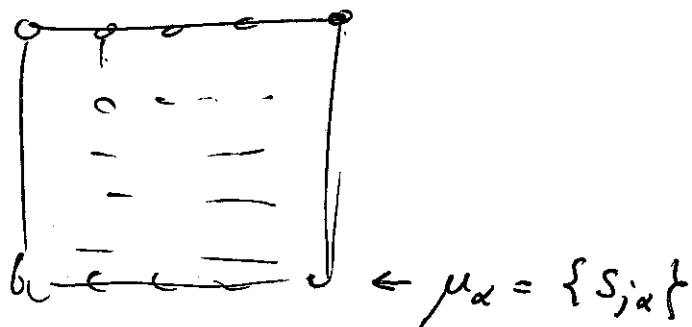


2d Ising model.

①



$$Z = \sum_{\{\mu_\alpha\}} e^{-\beta \sum_x E(\mu_x, \mu_{x+1}) - \sum_\alpha \beta E(\mu_\alpha)}$$

$$E_{\mu_\alpha} = - \sum_j s_j^\alpha s_{j+1}^\alpha$$

$$E_{\mu_1 \mu_2} = - \sum_j s_j^\alpha s_j^{\alpha-1}$$

$$\langle \mu | P | \mu' \rangle = e^{-\beta E(\mu, \mu') - \beta E(\mu)}$$

$$Z = \text{Tr } P^N = \lambda_1^N + \dots + \lambda_{2^N}^N$$

Assume  $\lambda_i \geq 0$        $\lambda_{\max}^N \leq Z \leq 2^N \lambda_{\max}^N$

$$N \ln \lambda_{\max} \leq \ln Z \leq N \ln 2 + N \ln \lambda_{\max}$$

$$\ln \lambda_{\max} \leq \frac{1}{N} \ln Z \leq \ln 2 + \ln \lambda_{\max}$$

but we have  $N^2$  sites

$$-\frac{\beta A}{N^2} = -\beta a$$

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$$\frac{1}{N} \ln \lambda_{\max} \leq -\beta a \leq \frac{1}{N} \ln 2 + \frac{1}{N} \ln \lambda_{\max}$$

$$N \rightarrow \infty \quad -\beta a = \lim_{N \rightarrow \infty} \frac{1}{N} \ln \lambda_{\max}$$

$$\langle s_1 \dots s_N | P | s'_1 \dots s'_N \rangle = \prod_{j=1}^N e^{\beta s_j} e^{\beta s_j s_{j+1}} e^{\beta s_j s'_j}$$

$$\langle s_1 \dots s_N | V_1' | s'_1 \dots s'_N \rangle = \prod_{j=1}^N e^{\beta s_j s'_j}$$

$$\langle s_1 \dots s_N | V_2 | s'_1 \dots s'_N \rangle = \delta_{s_1 s'_1} \dots \delta_{s_N s'_N} \prod_{j=1}^N e^{\beta s_j s_{j+1}}$$

$$\langle s_1 \dots s_N | V_3 | s'_1 \dots s'_N \rangle = \delta_{s_1 s'_1} \dots \delta_{s_N s'_N} \prod_{j=1}^N e^{\beta t s_j}$$

$$P = V_3 V_2 V_1'$$

$$V_1' = A \otimes \dots \otimes A$$

$$A = e^{\beta s_j s'_j} = \begin{pmatrix} e^{\beta} & e^{-\beta} \\ e^{-\beta} & e^{\beta} \end{pmatrix}$$

$$= e^{\beta} \mathbb{1} + e^{-\beta} \sigma_x$$

$$e^{\beta \sigma_x} = \cosh \theta + \sigma_x \sinh \theta = \cosh \theta (1 + \tanh \theta \sigma_x) = e^{\beta} (1 + e^{-2\beta} \sigma_x)$$

$$\tanh \theta = e^{-2\beta}$$

$$1 - \tanh^2 \theta = 1 - \frac{\sinh^2 \theta}{\cosh^2 \theta} = 1 - e^{-4\beta}$$

$$\cosh^2 \theta = \frac{1}{1 - e^{-4\beta}} = \frac{e^{2\beta}}{e^{2\beta} - e^{-2\beta}} = \frac{e^{2\beta}}{2 \sinh 2\beta}$$

$$e^{\beta \sigma_i} = \frac{e^{\beta}}{\sqrt{2 \sinh 2\beta}} (1 + e^{-2\beta} \sigma_i) = \frac{A}{\sqrt{2 \sinh 2\beta}}$$

$$A = \sqrt{2 \sinh 2\beta} e^{\beta \sigma_i}$$

$$V'_1 = (2 \sinh 2\beta)^{N/2} e^{\beta \sigma_1} \otimes \dots \otimes e^{\beta \sigma_1}$$

$$\sigma_1^\alpha = 1 \otimes \dots \otimes \sigma_1^\alpha \otimes \dots \otimes 1$$

$$V'_1 = (2 \sinh 2\beta)^{N/2} \prod_{\alpha=1}^N e^{\beta \sigma_1^\alpha} = (2 \sinh 2\beta)^{N/2} e^{\beta \sum_{\alpha=1}^N \sigma_1^\alpha}$$

commute

$$V_1 = e^{\beta \sum_{\alpha=1}^N \sigma_1^\alpha}$$

$$\tanh \theta = e^{-2\beta}$$

$$P = (2 \sinh 2\beta)^{N/2} V_3 V_2 V_1$$

$$V_3 = A_3 \otimes \dots \otimes A_3$$

$$A_3 = \begin{pmatrix} e^{\beta H} & 0 \\ 0 & e^{-\beta H} \end{pmatrix} = e^{\beta H \sigma_3}$$

$$\langle s_i, \dots, s_n | V_3 | s'_1, \dots, s'_n \rangle = A_{3, s_1 s'_1} \dots A_{3, s_n s'_n} = e^{\beta H \sum \sigma_3^{(i)}}$$

$$\langle s_i, \dots, s_n | V_2 | s'_1, \dots, s'_n \rangle = \delta_{s_1 s'_1} \dots \delta_{s_n s'_n} \prod_{j=1}^N e^{\beta S_j s_{j+1}}$$

$$\prod_{\alpha=1}^N e^{\beta \sigma_3^\alpha \sigma_3^{\alpha+1}}$$

$$e^{\beta \sigma_3^\alpha \sigma_3^{\alpha+1}}$$

$$\sigma_3^\alpha \sigma_3^{\alpha+1} = 1 \otimes \dots \otimes \overset{\alpha}{\sigma_3} \otimes \overset{\alpha+1}{\sigma_3} \otimes \dots \leftarrow \text{diagonal!}$$

$$\langle s_i, \dots, s_n | e^{\beta \sigma_3^\alpha \sigma_3^{\alpha+1}} | s'_1, \dots, s'_n \rangle = \delta_{s_1 s'_1} \dots \delta_{s_n s'_n} \cdot e^{\beta S_j s_{j+1}}$$

$$V_2 = \prod_{\alpha=1}^N e^{\beta \sigma_3^\alpha \sigma_3^{\alpha+1}}$$

$$V_1 = e^{\theta \sum_{\alpha=1}^N \sigma_1^\alpha}$$

$\tanh \theta = e^{-2/\beta}$

If  $\beta=0$

$$P = (2 \sinh 2\beta)^{N/2} V_3 V_2 V_1$$

$$V_3 = e^{\beta H \sum \sigma_3^{(i)}}$$

Fermionic representation

$$|0\rangle |1\rangle \dots |1\rangle$$

$$c_j^\dagger |0\rangle = |1\rangle$$

$$c_j |1\rangle = |0\rangle$$

$$\{c_i, c_j^\dagger\} = \delta_{ij}$$

anti-commute at different sites.

1-site

$$c = \begin{pmatrix} |0\rangle & |1\rangle \\ \langle 0| & \langle 1| \end{pmatrix}$$

$$c = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

but at different sites they would ~~not~~ commute.

take  $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$  ( $\frac{1 \pm \sigma_3}{2}$  = fermion numbers)

$$\sigma_3 \cdot c^\dagger + c^\dagger \cdot \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ -1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} = 0$$

$$\sigma_3 \cdot c + c \cdot \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & -1 \\ 0 & 0 \end{pmatrix} = 0$$

$$C_j = \sigma_3 \otimes \dots \otimes \sigma_3 \otimes c \otimes \mathbb{1} \dots \otimes \mathbb{1}$$

$$\bar{C}_j = \sigma_3 \otimes \dots \otimes \sigma_3 \otimes c^\dagger \otimes \mathbb{1} \dots \otimes \mathbb{1}$$

↑ acting on each site.

$$V_1 = e^{\sum_{\alpha=1}^N \sigma_1^\alpha}$$

rotation  
(change of basis)

$$e^{\sum_{\alpha=1}^N \sigma_3^\alpha}$$

$$V_2 = e^{\sum_{\alpha} \sigma_3^\alpha \sigma_3^{\alpha+1}}$$

→

$$e^{\sum_{\alpha=1}^{N-1} \sigma_1^\alpha \sigma_1^{\alpha+1}}$$

↑ consider open b.c. instead of periodic.

$$c c^\dagger = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \quad c^\dagger c = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

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$$-c^\dagger c + c c^\dagger = \sigma_3$$

$$c_j^\dagger c_j - c_j c_j^\dagger = 1 \otimes \dots \otimes \sigma_3 \otimes \dots \otimes 1 = \sigma_3^{(j)}$$

$$V_1 = e^{\sum_j (c_j^\dagger c_j - c_j c_j^\dagger)}$$

$$i(c^\dagger - c) = \sigma_2 \quad c + c^\dagger = \sigma_1 \quad \sigma_3 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma_1 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$c_{j+1} + c_{j+1}^\dagger = \sigma_3 \otimes \dots \otimes \sigma_3 \otimes \sigma_1 \otimes \dots \otimes 1$$

$$i(c_j^\dagger - c_j) = \sigma_3 \otimes \dots \otimes \sigma_2 \otimes \dots \otimes 1 \otimes \dots \otimes 1$$

$$i(c_j^\dagger - c_j)(c_{j+1} + c_{j+1}^\dagger) = 1 \otimes \dots \otimes 1 \otimes i\sigma_2 \otimes \sigma_1 \otimes \dots \otimes 1 = i\sigma_2 \otimes \sigma_1$$

$$= i\sigma_1^{(j)} \cdot \sigma_1^{(j+1)}$$

$$V_2 = e^{\beta \sum_{j=1}^{N-1} (c_j^\dagger - c_j)(c_{j+1} + c_{j+1}^\dagger)} + \beta \left( \prod_{j=1}^N \sigma_3 \right) (c_N^\dagger - c_N)(c_1 + c_1^\dagger)$$

$$\sigma_2 \cdot \sigma_3 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} = i\sigma_1$$

$$\begin{aligned} \sigma_3 \otimes \sigma_3 &= \sigma_3 \otimes \sigma_3 \\ \sigma_2 \otimes 1 &= \sigma_2 \otimes 1 \\ -i\sigma_2 \otimes \sigma_3 &= \sigma_2 \otimes \sigma_2 \\ \sigma_3 \otimes \sigma_3 &= \sigma_3 \otimes \sigma_3 \end{aligned}$$

$$\begin{aligned} \beta \cdot \sigma_1 &= -i\sigma_2 \\ (i\sigma_2) \otimes 1 &= \sigma_1 \otimes \sigma_1 \\ = i\sigma_1 \otimes 1 &= \sigma_1 \otimes \sigma_1 \end{aligned}$$

Fermionic representation

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1 fermion in each site

$$P = (2 \sinh 2\beta)^{N/2} V_2 V_1$$

$$V_2 = e^{\beta \sum_i (c_{j-} - c_{j+}^{\dagger})(c_{j+1} + c_{j+1}^{\dagger})}$$

$$V_1 = e^{\theta \sum_i (c_{j-}^{\dagger} \epsilon_j - c_{j-} c_{j+}^{\dagger})}$$

Define  $\{c_p, c_{p'}^{\dagger}\} = \delta(p-p')$   $-\pi \leq p \leq \pi$

$$c_j = \int_{-\pi}^{\pi} \frac{dp}{\sqrt{2\pi}} e^{ipj} c_p \quad ; \quad c_j^{\dagger} = \int_{-\pi}^{\pi} \frac{dp}{\sqrt{2\pi}} e^{-ipj} c_p^{\dagger}$$

$$\{c_j, c_{j'}^{\dagger}\} = \int_{-\pi}^{\pi} \frac{dp dp'}{2\pi} e^{i(pj - p'j')} \delta(p-p') = \int_{-\pi}^{\pi} \frac{dp}{2\pi} e^{ip(j-j')}$$

$\begin{matrix} \nearrow j=j' & \rightarrow 1 \checkmark \\ \searrow j \neq j' & \rightarrow 0 \end{matrix}$

$$\frac{e^{ip(j-j')}}{2\pi i(j-j')} \Big|_{-\pi}^{\pi} = 0 \quad (j-j' \text{ integer})$$

$$N \int_{-\pi}^{\pi} \frac{dp}{2\pi} = 1 \quad \text{density of states.}$$

Also 
$$c_j = \sum_{n=0}^{N-1} \frac{e^{2\pi i n j}}{\sqrt{N}} c_n$$

$$\sum_{j=-\infty}^{\infty} c_j^\dagger c_j - c_j^\dagger c_j^\dagger = \int_{-\pi}^{\pi} \frac{dp dp'}{2\pi} \sum_j e^{-ipj + ip'j} c c_p^\dagger c_{p'} - c_p c_p^\dagger \quad (8)$$

periodic

$$= \int_{-\pi}^{\pi} dp (c_p^\dagger c_p - c_p c_p^\dagger)$$

$$\sum_j (c_j - c_j^\dagger) (c_{j+1} + c_{j+1}^\dagger) =$$

$$= \int_{-\pi}^{\pi} \frac{dp dp'}{2\pi} (e^{ipj} c_p - e^{-ipj} c_p^\dagger) (e^{ip'(j+1)} c_{p'} + e^{-ip'(j+1)} c_{p'}^\dagger)$$

$$= \int_{-\pi}^{\pi} dp (e^{-ip} c_p c_{-p} + e^{-ip} c_p c_p^\dagger - e^{+ip} c_p^\dagger c_{+p} - c_p^\dagger c_p^\dagger)$$

We can look at each  $p$  separately.  $(p, -p)$

$|00\rangle$   $|01\rangle$   $|10\rangle$   $|11\rangle$   
 $p-p$   $p-p$   $p-p$   $p-p$

fermionic parity.  
↑

But each term conserves even or odd # of fermions  
 since it conserves # of odd or subtracts 2.



$$V_1 = e^{i\theta \int_0^\pi dp (c_p^\dagger c_p - c_p c_p^\dagger + \frac{c_p^\dagger}{-p} c_{-p} - c_{-p} c_p^\dagger)} \quad (9)$$

$$V_2 = e^{i\beta \int_0^\pi dp (e^{-ip} c_p c_{-p} + e^{ip} c_{-p} c_p + e^{-ip} c_p c_p^\dagger + e^{ip} c_{-p} c_{-p}^\dagger - e^{ip} c_p^\dagger c_p - e^{-ip} c_{-p}^\dagger c_{-p} - e^{ip} c_p^\dagger c_{-p}^\dagger - e^{-ip} c_{-p}^\dagger c_p^\dagger)}$$

Consider basis  $|00\rangle |11\rangle |10\rangle |11\rangle$  for  $p, -p$  states.

$$\begin{matrix} \langle 00| \\ \langle 11| \\ \langle 01| \\ \langle 10| \end{matrix} \begin{pmatrix} |00\rangle & |11\rangle & |10\rangle & |11\rangle \end{pmatrix}$$

$$C_p = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \quad C_p^\dagger = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix} \quad C_{-p} = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix} \quad C_{-p}^\dagger = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$V_1 = e^{-2\theta \sum_p \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}}$$

$$V_2 = e^{2\beta \sum_p \begin{pmatrix} cp & isp & 0 & 0 \\ -isp & -cp & 0 & 0 \\ 0 & 0 & c & c \\ 0 & c & c & c \end{pmatrix}}$$

looking at  $(2 \times 2)$  part (the other 2 eigenvalues are 0).

$$V_1 = e^{-2\theta \sum_p \sigma_3}$$

$$V_2 = e^{2\beta \sum_p (cp \sigma_3 - isp \sigma_2)}$$

$$e^{-2\theta \sigma_3} e^{2\beta (cp \sigma_3 - isp \sigma_2)} = (\cosh 2\theta - \sinh 2\theta \sigma_3) \cdot (\cosh \beta + \sinh \beta (cp \sigma_3 - isp \sigma_2))$$

$$= \left[ \cosh 2\theta \cosh \beta + \cosh 2\theta \sinh \beta (cp \sigma_3 - isp \sigma_2) - \sinh 2\theta \cosh \beta \sigma_3 - \sinh 2\theta \sinh \beta (cp \sigma_3 - isp \sigma_2) \right] = \cosh \mu + \sinh \mu (\hat{n} \cdot \sigma)$$

squares to 1  
eigenvalues  $\pm 1$

$$e^{-2\theta\sigma_3} e^{2\beta(cp\sigma_3 - sp\sigma_2)} = e^{\mu(\hat{n}\cdot\sigma)} \quad (10)$$

eigenvalues  $\pm\mu, \pm\mu, \sigma, 0$  symmetric  $\beta \leftrightarrow \tilde{\beta}$ .

$$\sum_P \mu(\theta) = \frac{N}{2\pi} \int_0^\pi d\phi \operatorname{arccosh}(\operatorname{ch} 2\theta \operatorname{ch} 2\beta - c\phi)$$

$\operatorname{sh} 2\theta \operatorname{sh} 2\beta = 1$   $\operatorname{ch} 2\theta = \frac{\operatorname{ch}^2 2\beta}{\operatorname{sh} 2\beta}$  if  $\mu(\theta) = 1$   
we get  $N \checkmark$ .

$$\operatorname{Tr} P^N = e^{\frac{N^2}{2\pi} \int_0^\pi d\phi \operatorname{arccosh} \left( \frac{\operatorname{ch}^2 2\beta}{\operatorname{sh} 2\beta} - c\phi \right)} \cdot (2\operatorname{sh} 2\beta)^{N^2/2}$$

$$e^{-\beta A} = \operatorname{Tr} P^N$$

$$-\beta a = -\beta \frac{A}{N^2} = \frac{1}{2} \ln(2\operatorname{sh} 2\beta) + \frac{1}{2\pi} \int_0^\pi d\phi \operatorname{arccosh} \left( \frac{\operatorname{ch}^2 2\beta}{\operatorname{sh} 2\beta} - \cos \phi \right)$$

$$\beta a_{\text{asy}} = -\frac{1}{2} \ln(2\operatorname{sh} 2\beta) - \frac{1}{2\pi} \int_0^\pi d\phi \operatorname{arccosh} \left( \frac{\operatorname{ch}^2 2\beta}{\operatorname{sh} 2\beta} - \cos \phi \right)$$

$\beta a_{\text{asy}}$  same  $\checkmark$

$$= -\ln(2\operatorname{ch} 2\beta) - \frac{1}{2\pi} \int_0^\pi d\phi \ln \left( \frac{1}{2} (1 + \sqrt{1 - k^2 \sin^2 \phi}) \right)$$

both  $\rightarrow$

$$k^2 = \frac{2\operatorname{sh} 2\beta}{\operatorname{ch}^2 2\beta}$$