

Spring, 2022

ME 597 – Solid Mechanics II

Lecture 6 Thermodynamics

KEEP A MASK WITH
YOU AT ALL TIMES



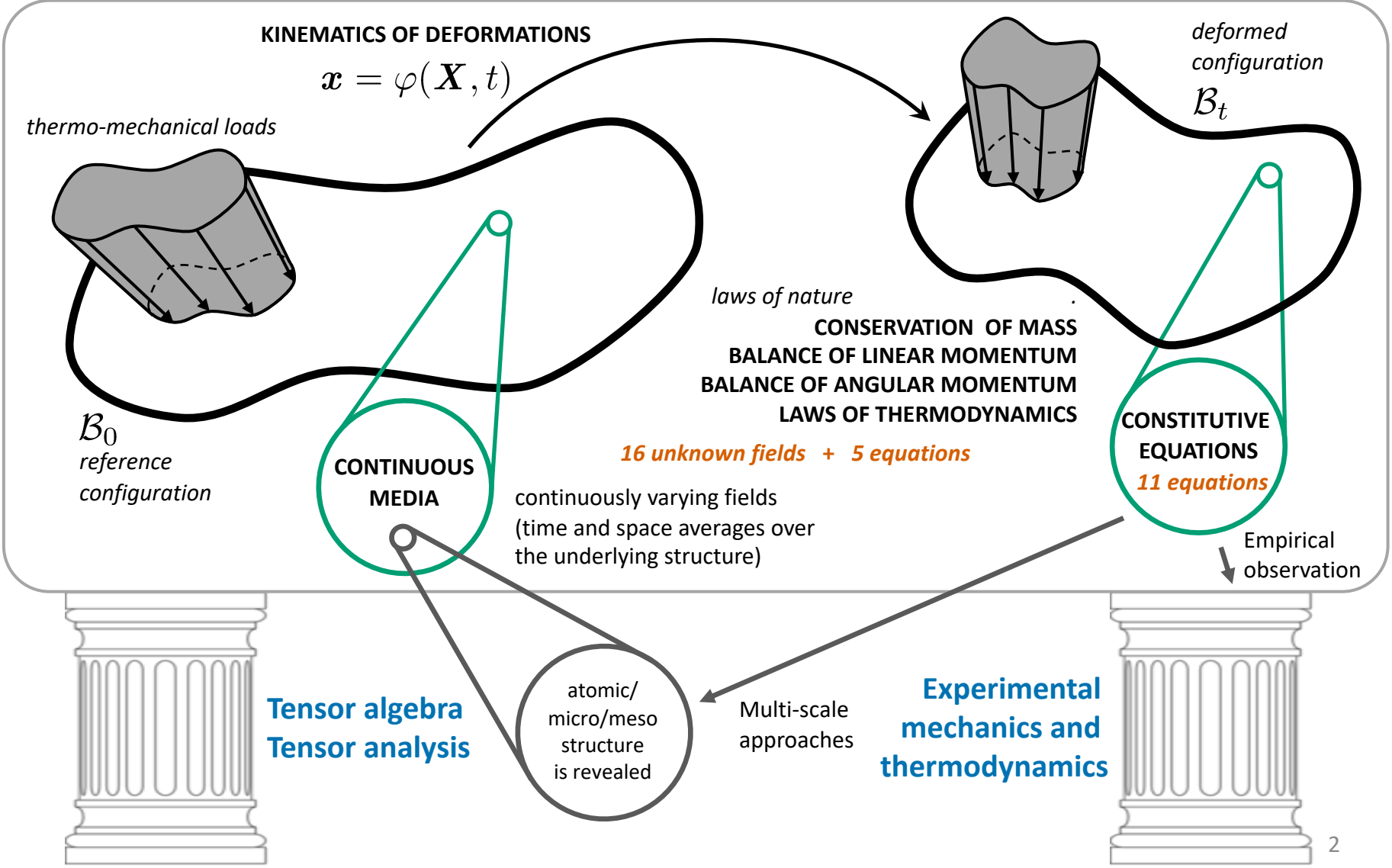
**PROTECT
PURDUE**



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Lecture 6 – Thermodynamics

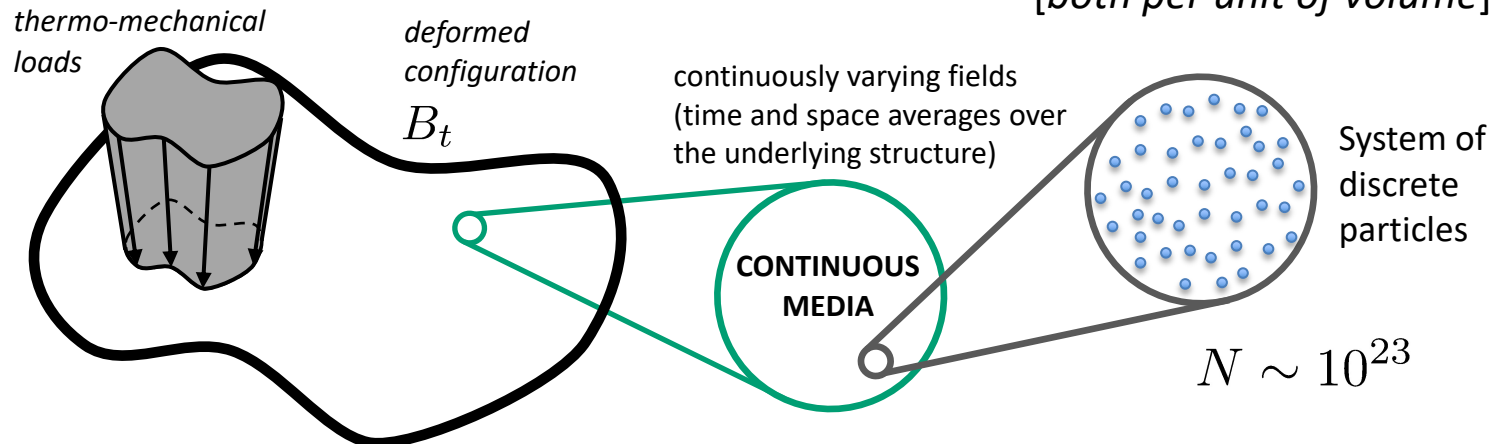


Macroscopic observable quantities

- Microscopic kinematics are described by a $6N$ -dimensional vector space (phase space) with all particle positions and momenta

$$\mathbf{y} = \{ \mathbf{r}^1, \dots, \mathbf{r}^N, m^1 \dot{\mathbf{r}}^1, \dots, m^N \dot{\mathbf{r}}^N \}$$

- Macroscopic kinematic quantities (e.g., total volume)
- Macroscopic non-kinematic quantities (e.g., total number of particles, total mass)
- Macroscopic fields (e.g., density, Lagrangian strain tensor)
[both per unit of volume]



Thermodynamic equilibrium

After an external perturbation (of finite duration in time), all systems tend to evolve to a quiescent and spatially homogeneous (at the macroscopic length scale) terminal state where the system's macroscopic observables have constant limiting values.

State variables

Macroscopic observables that are well defined and single-valued when the system is at state of thermodynamic equilibrium.

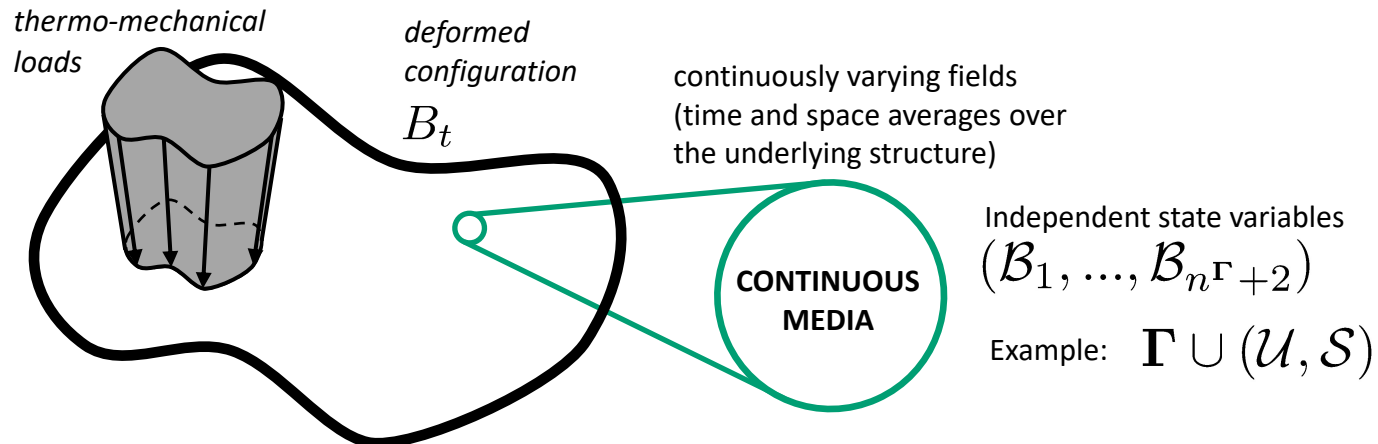
Example: the thermodynamic equilibrium of a gas is characterized by the number of particles, the volume, the pressure and the temperature (i.e., 4 state variables).

Some state variables are kinematic state variables (such as volume and, for solids, strain). A change in a kinematic state variable results in a *thermodynamic tension*.

Thermodynamics

Laws of thermodynamics

- Zeroth law of thermodynamics $T \geq 0$ (...concept of absolute temperature)
- First law of thermodynamics \mathcal{U} (...statement of conservation of energy)
(...it leads to the idea of internal energy)
- Second law of thermodynamics (...directionality of thermodynamic processes)
for an isolated system $\Delta \mathcal{S} \geq 0$ (...it leads of the idea of entropy)



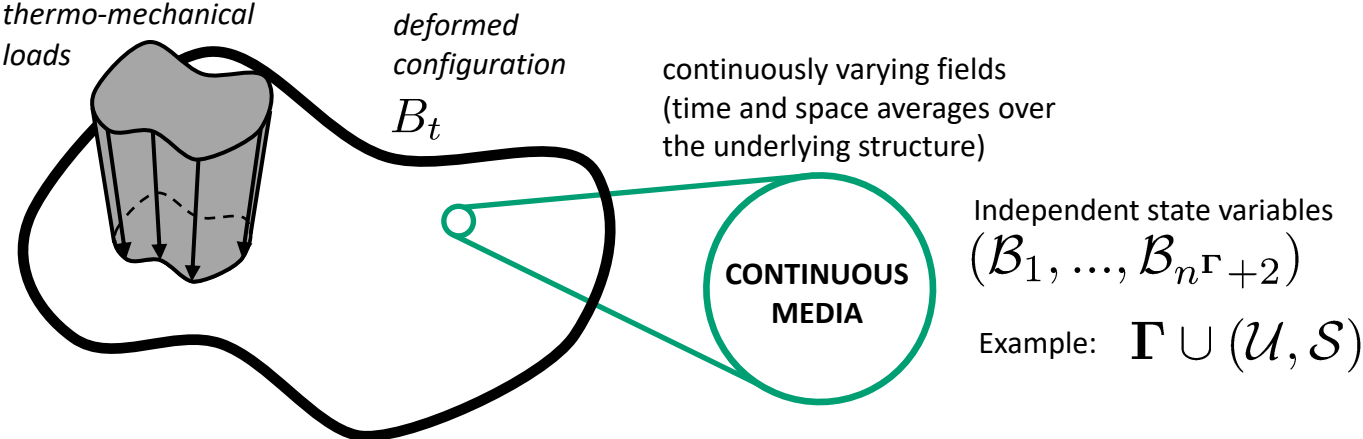
Thermodynamics

(back to ...) State variables

Any system in thermodynamic equilibrium is fully characterized by $n^\Gamma + 2$ independent state variables, where n^Γ is the number of independent kinematic state variables $\Gamma = (\Gamma_1, \dots, \Gamma_{n^\Gamma})$ that characterize the system.

$\mathcal{B} = (\mathcal{B}_1, \dots, \mathcal{B}_{n^\Gamma+2}, \dots)$ $\mathcal{B}_{n^\Gamma+2+j} = f_j(\mathcal{U}, \mathcal{S}, \Gamma_1, \dots, \Gamma_{n^\Gamma}) \quad j = 1, 2, \dots$
 All state variables Equations of state

Gas:
 Solid: DIY



Thermodynamics

Continuum thermodynamics

- Postulate of local thermodynamic equilibrium

The local and instantaneous relations between thermodynamic quantities in a system out of equilibrium are the same as for a uniform system in equilibrium.

$$T \equiv T(\mathbf{x}, t)$$

$$p \equiv p(\mathbf{x}, t)$$

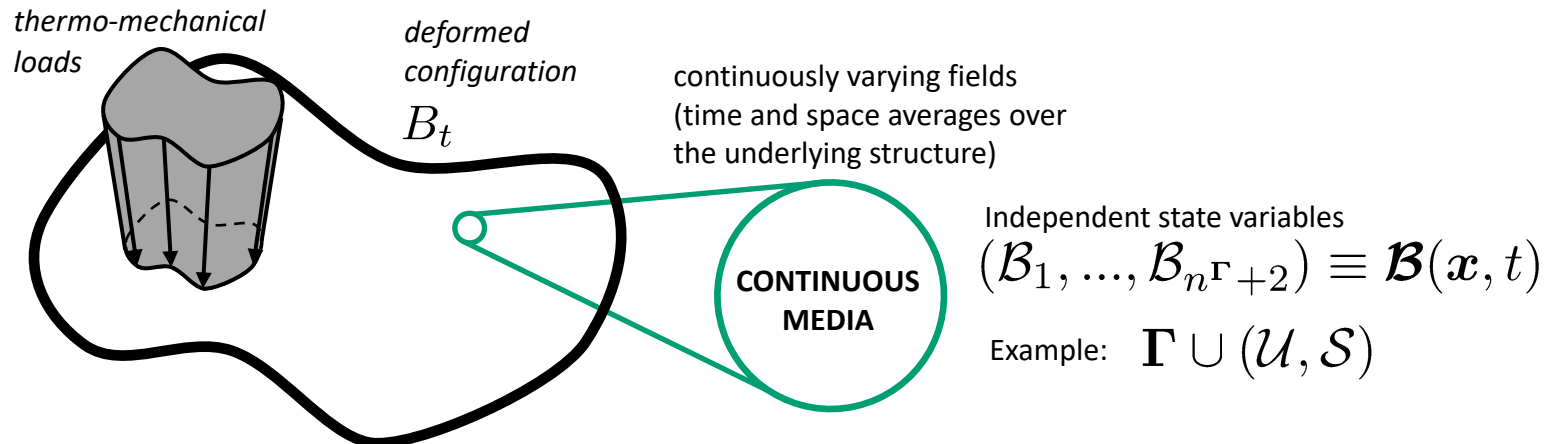
$$\rho \equiv \rho(\mathbf{x}, t)$$

$$u \equiv u(\mathbf{x}, t)$$

$$s \equiv s(\mathbf{x}, t)$$

specific internal energy
(internal energy per unit mass)

specific entropy
(entropy per unit mass)



Thermodynamics

Continuum thermodynamics

- First law of thermodynamics

$$\rho \dot{u} = \sigma_{ij} d_{ij} + \rho r - q_{i,i} \iff$$

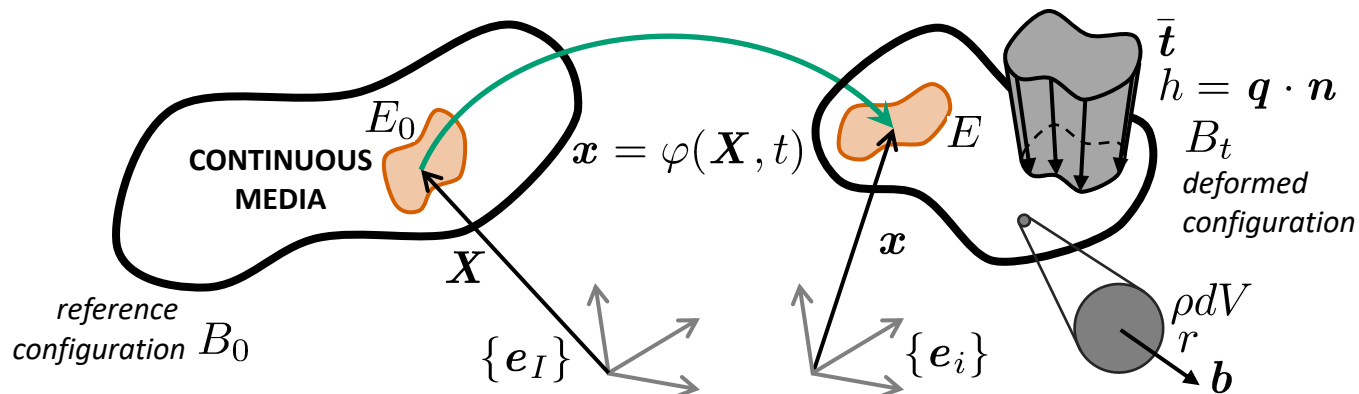
$$\rho \dot{u} = \boldsymbol{\sigma} : \mathbf{d} + \rho r - \operatorname{div} \mathbf{q} \quad \forall \mathbf{x} \in B_t$$

spatial form of the conservation of energy

- Second law of thermodynamics

$$\dot{s} \geq \dot{s}^{\text{ext}} = \frac{r}{T} - \frac{1}{\rho} \operatorname{div} \frac{\mathbf{q}}{T} \quad \forall \mathbf{x} \in B_t$$

Clausius-Duhem inequality



Continuum mechanics – Laws of nature

Summary

$$J\rho = \rho_0 \quad \forall \mathbf{X} \in B_0 \quad \text{conservation of mass} \quad (1 \text{ equation})$$

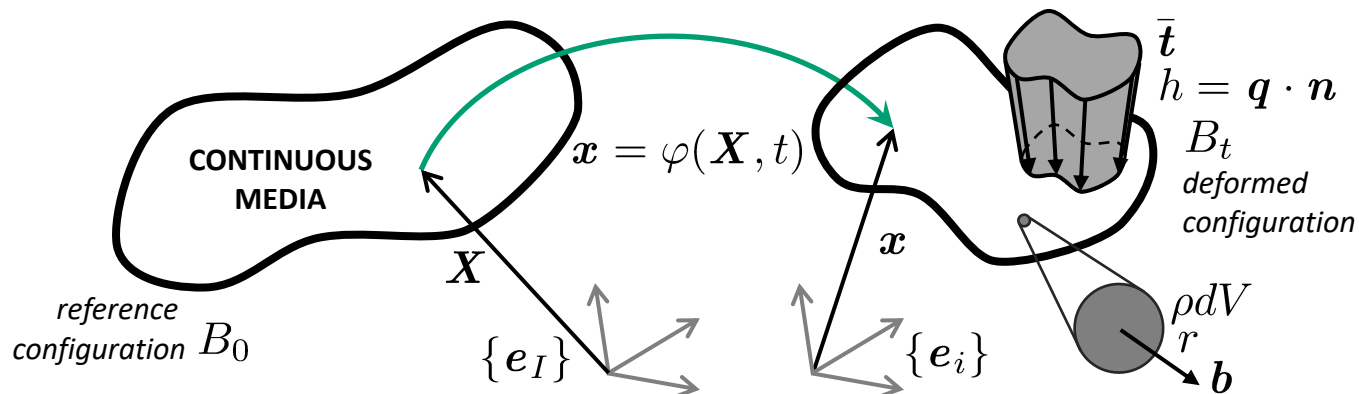
$$\operatorname{div} \boldsymbol{\sigma} + \rho \mathbf{b} = \rho \mathbf{a} \quad \forall \mathbf{x} \in B \quad \text{balance of linear momentum} \quad (3 \text{ equations})$$

$$\boldsymbol{\sigma} = \boldsymbol{\sigma}^T \quad \forall \mathbf{x} \in B \quad \text{balance of angular momentum} \quad (\text{constraint})$$

$$\rho \dot{u} = \boldsymbol{\sigma} : \mathbf{d} + \rho r - \operatorname{div} \mathbf{q} \quad \forall \mathbf{x} \in B \quad \text{conservation of energy} \quad (1 \text{ equation})$$

$$\dot{s} \geq \frac{r}{T} - \frac{1}{\rho} \operatorname{div} \frac{\mathbf{q}}{T} \quad \forall \mathbf{x} \in B \quad \text{Clausius-Duhem inequality} \quad (\text{constraint})$$

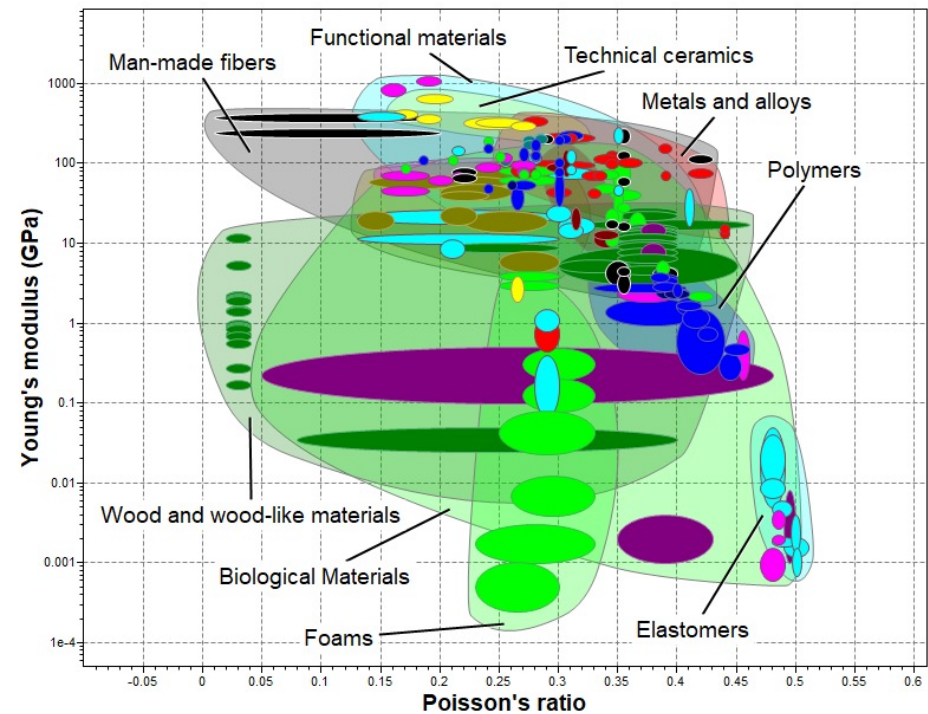
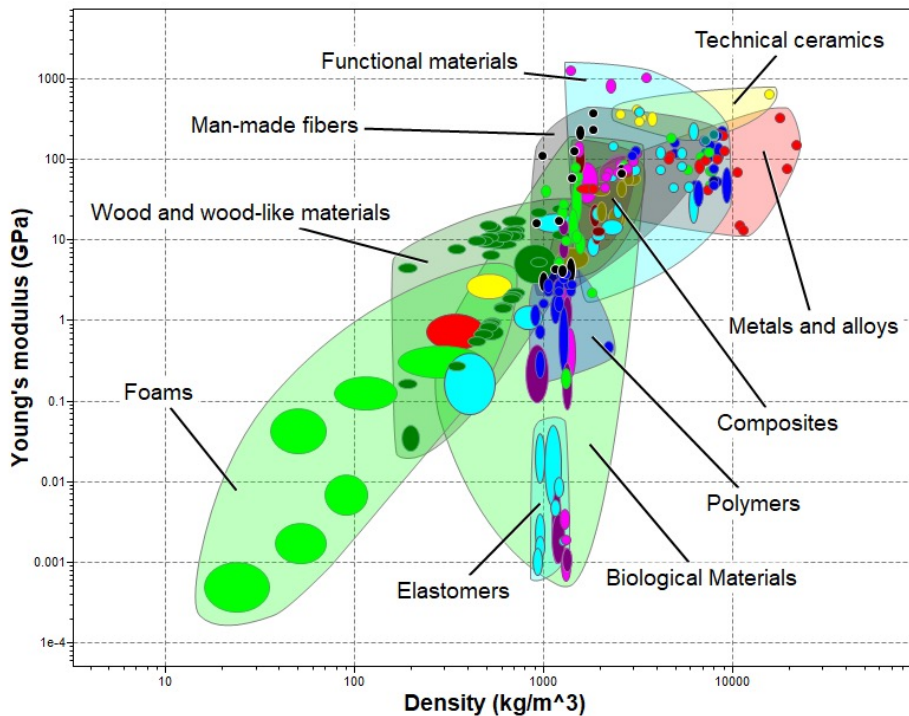
$$\rho, \mathbf{x}, \boldsymbol{\sigma}, \mathbf{q}, u, s, T \quad (16 \text{ unknowns})$$



ME 597 – Solid Mechanics II (Ashby plots)

Young's modulus – Poisson's ratio – Density

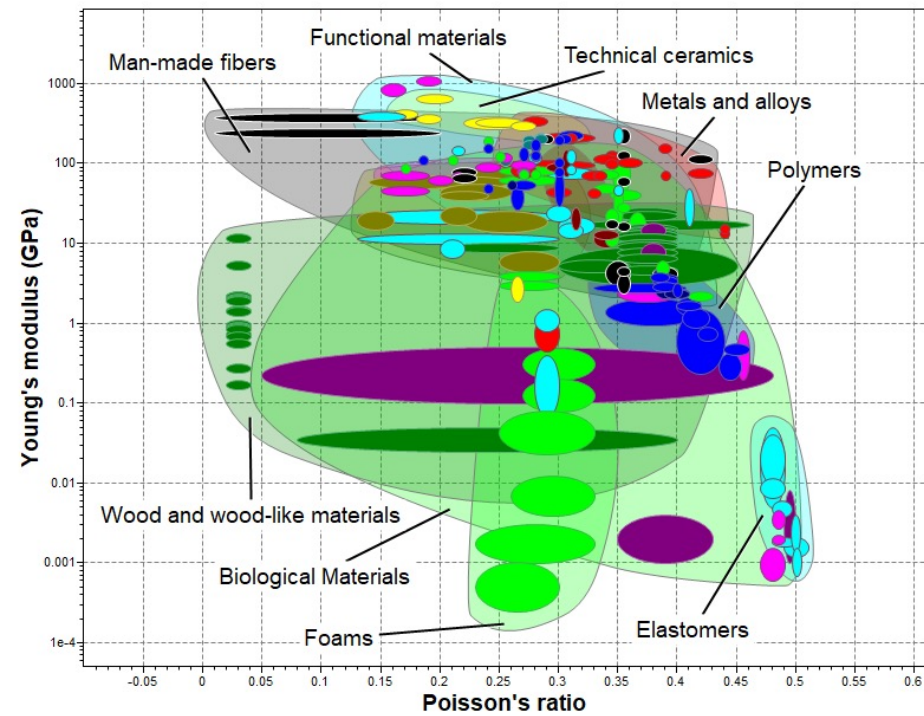
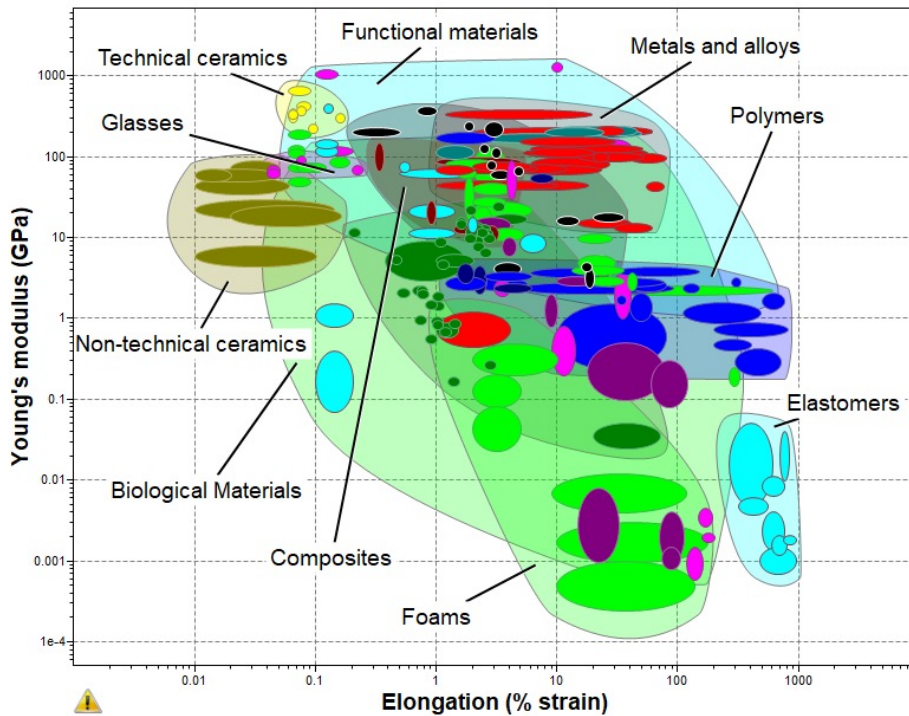
- Flexible materials are typically light, and stiff materials are typically heavy.
- Polymers and elastomer are almost incompressible solids.
- Stiff lightweight materials are hard to find.



ME 597 – Solid Mechanics II (Ashby plots)

Young's modulus – Poisson's ratio – Elongation at failure

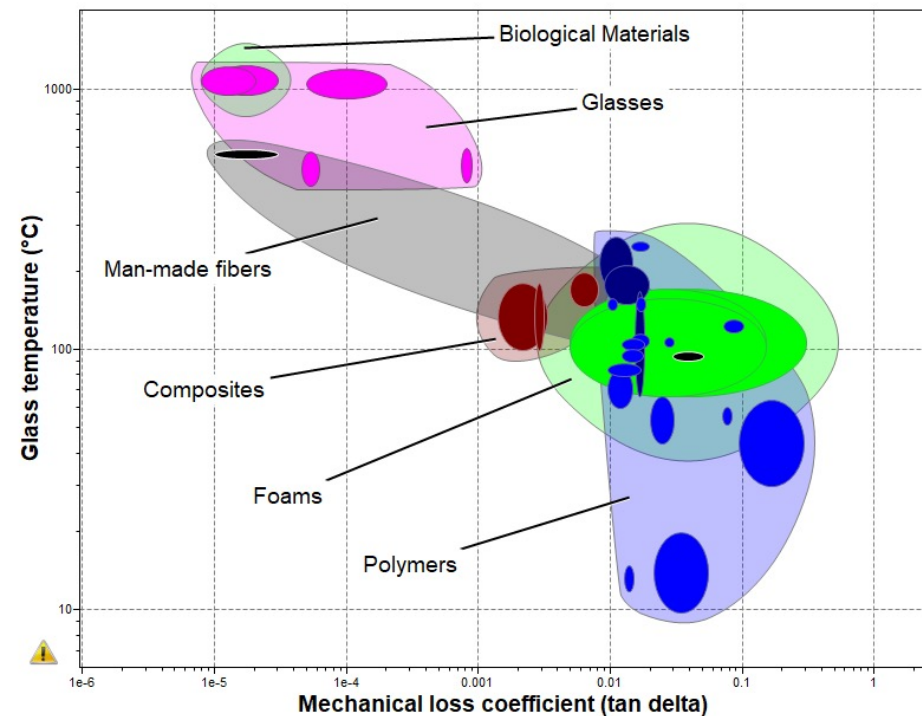
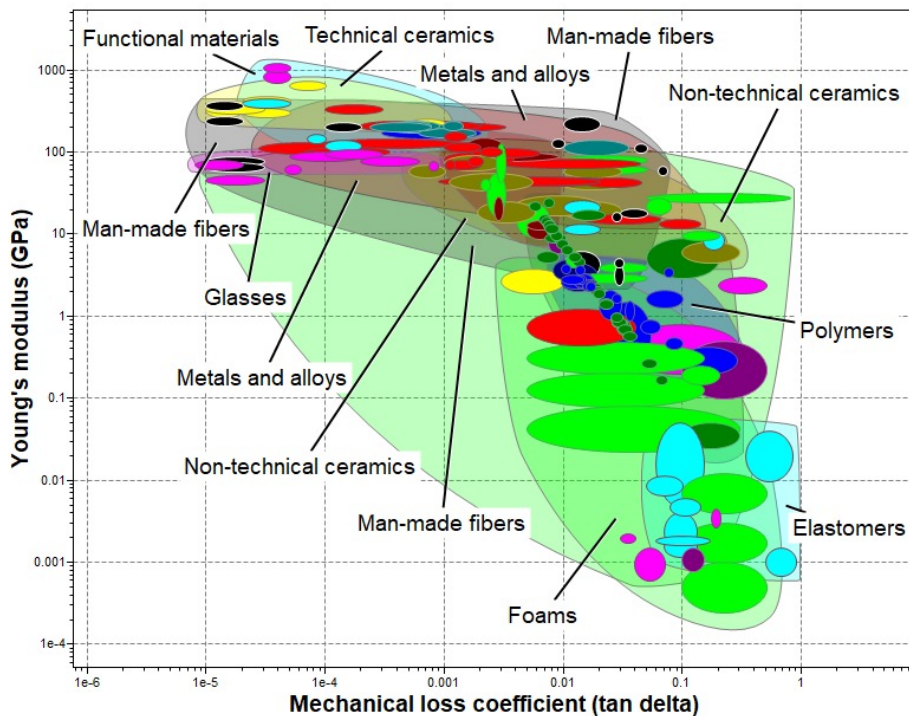
- Elastomers are polymers with viscoelasticity.
- Polymers and elastomers can exhibit elongations over 200%!
- Polymers and elastomer are almost incompressible solids.



ME 597 – Solid Mechanics II (Ashby plots)

Young's modulus – Mechanical loss – Glass temperature

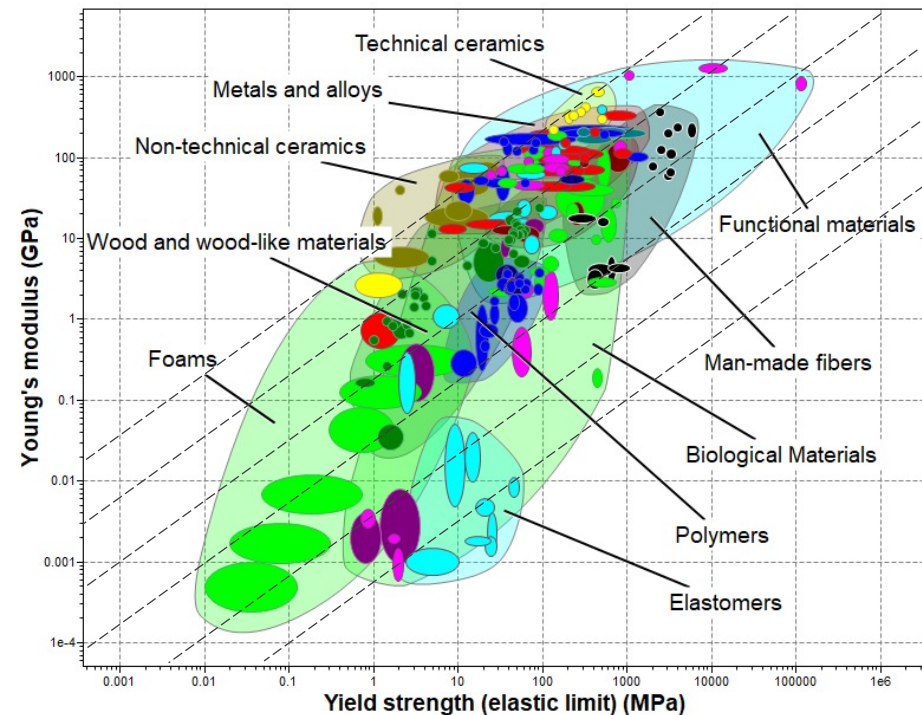
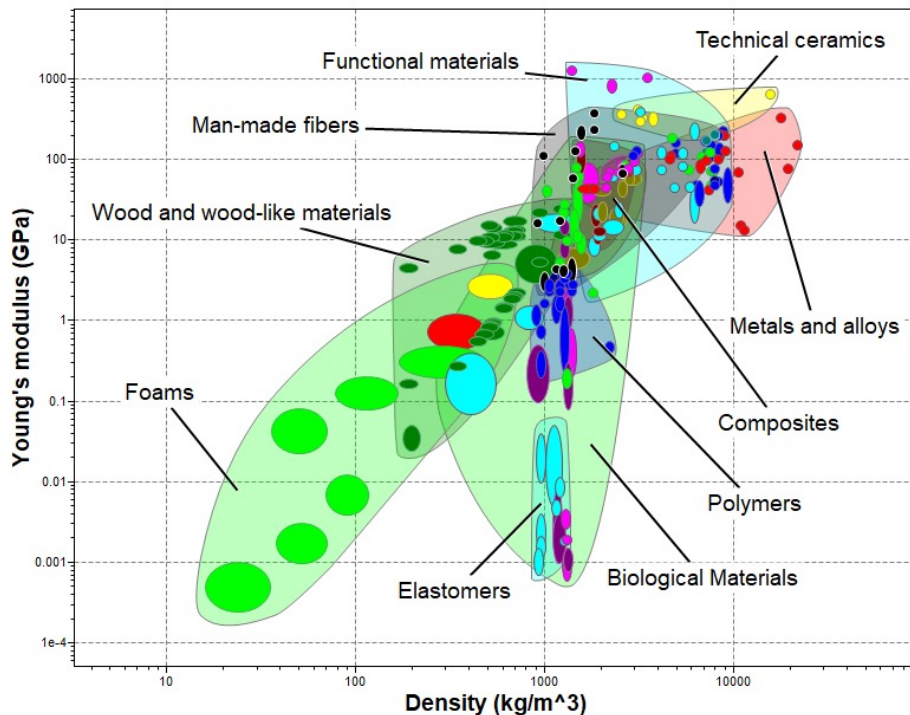
- Reversible transition from hard/brittle 'glassy' state into viscous state.
- Materials with good damping behavior (i.e., high vibrational energy dissipation) are typically soft and have low glass-transition temperatures.



ME 597 – Solid Mechanics II (Ashby plots)

Young's modulus – Yield stress– Density

- Irreversible transition from elastic to plastic behavior.
- Contours of constant yield strain are a family of parallel lines on the Young's modulus vs. yield strength plot.



Lecture 6 – Thermodynamics

Any questions?