Lecture 21 Structural elements: beams, plates, shells





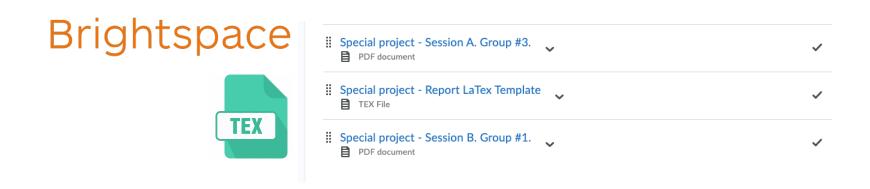
Mechanical Engineering Instructor: Prof. Marcial Gonzalez

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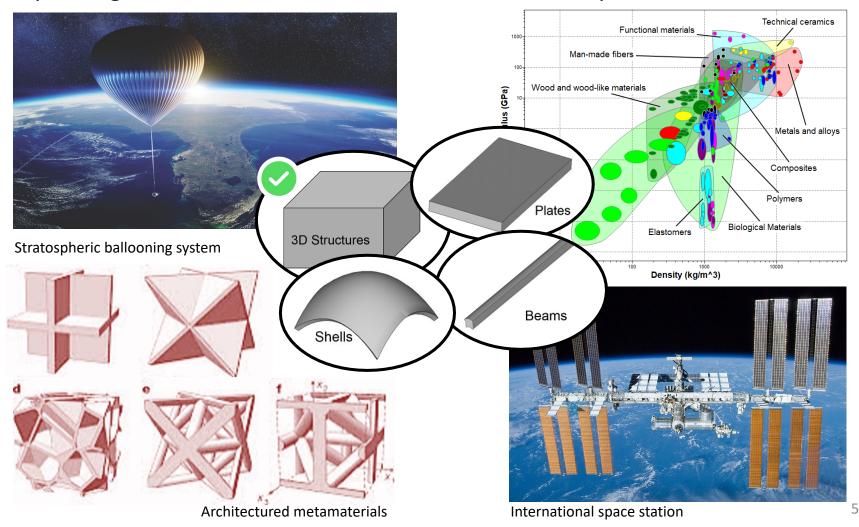
Announcements

Guidelines for special project. Final report:

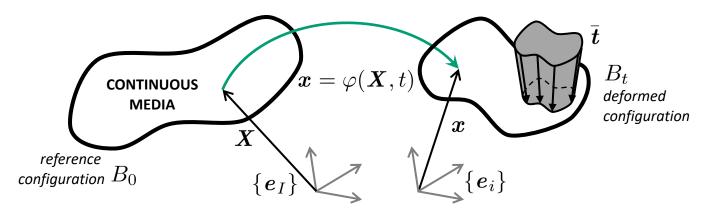
- Monday May 2nd at noon
- Technical report in LaTex, maximum of 4 pages (upload to Brightspace)
- Rewrite equations using the notation and nomenclature adopted in the class
- Grading: 5% of final grade (17% of project grade)



Formulation of structural elements (beams, plates, shells) as the analytical upscaling of continuum solids under kinematic assumptions.



Formulation of structural elements (beams, plates, shells) as the analytical upscaling of continuum solids under kinematic assumptions.

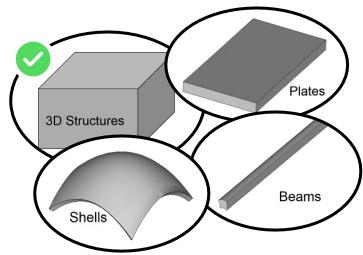


Linearized kinematic – Small trains

- The linearization is evaluated in the undeformed configuration (i.e., $m{X} o m{X} + m{u}(m{X})$ and $m{F} = m{I}, \, \nabla_0 m{u} = \nabla m{u}$): $\langle \nabla_\varphi m{E}; m{u} \rangle = \frac{1}{2} [\nabla m{u} + (\nabla m{u})^T] = m{\epsilon} \qquad \begin{array}{c} \text{small-strain tensor} \\ \text{(employed in elasticity theory)} \end{array}$
- Simplified geometry
- Simplifying kinematic assumptions (i.e., simplified displacement field $m{u}=(u_1,u_2,u_3)$)

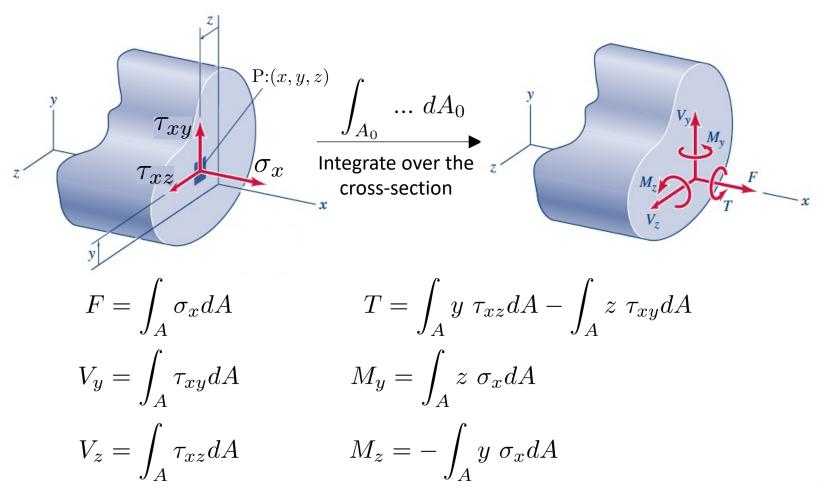
Simplified geometry

- Many everyday engineering applications utilize structural members such as rods, beams, cables, plates, and shells.
- Structural members can be idealized as one-dimensional (rods, beams, and cables) and two-dimensional (plates and shells) members.
- Assumption: two dimensions
 (in the case of beams) or
 one dimension (in the case of
 plates and shells) are significantly
 smaller than the other dimensions.



 Stress-strain behavior is upscaled to a relationship between internal resultants and kinematic variables (of the mid-plane).

Review (undergrad): force and moment resultants



Timoshenko beams

- Kinematic assumptions, with $ar{w} \ll h$

$$u_1(x, z) = \bar{u}(x) + z\theta_2(x)$$

$$u_2 = 0$$

$$u_3(x) = \bar{w}(x)$$

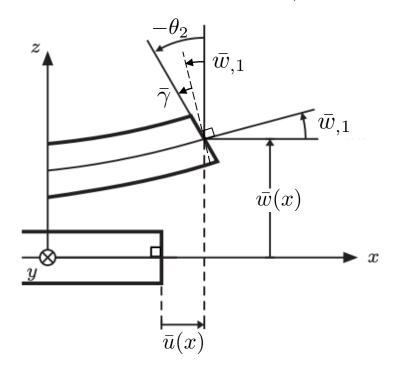
- Small-strain tensor, with $(\bar{w}_{,1})^2 \ll 1$

$$\epsilon_{ij} = \frac{1}{2}(u_{i,j} + u_{j,i})$$

$$\epsilon_{ij} = \begin{bmatrix} \bar{u}_{,1} + z\theta_{2,1} & 0 & (\theta_2 + \bar{w}_{,1})/2\\ 0 & 0 & 0\\ (\theta_2 + \bar{w}_{,1})/2 & 0 & 0 \end{bmatrix}$$

$$\epsilon_{ij} = \begin{bmatrix} \bar{\epsilon}_{11} + z\bar{\kappa}_{11} & 0 & \bar{\gamma}_{13}/2 \\ 0 & 0 & 0 \\ \bar{\gamma}_{13}/2 & 0 & 0 \end{bmatrix}$$

Simplified geometry: $b \ll L$, $h \ll L$



kinematic variables

 $ar{\epsilon}_{11}(x)$ $ar{\kappa}_{11}(x)$ $ar{\gamma}_{13}(x)$ extensional, bending, and shear components of beam strain

Timoshenko beams

- Small-strain tensor || Linear elasticity (generalized Hooke's law)

$$\epsilon_{ij} = \begin{bmatrix} \bar{\epsilon}_{11} + z\bar{\kappa}_{11} & 0 & \bar{\gamma}_{13}/2 \\ 0 & 0 & 0 \\ \bar{\gamma}_{13}/2 & 0 & 0 \end{bmatrix} \qquad \sigma_{ij} = \begin{bmatrix} E(\bar{\epsilon}_{11} + z\bar{\kappa}_{11}) & 0 & \frac{E}{2(1+\nu)}\bar{\gamma}_{13} \\ 0 & 0 & 0 \\ \frac{E}{2(1+\nu)}\bar{\gamma}_{13} & 0 & 0 \end{bmatrix}$$

... assuming $\sigma_{22} \ll \sigma_{11}$ and $\sigma_{33} \ll \sigma_{11}$ Since lateral surfaces are traction free, stresses in 22 and 33 directions must be zero on free surfaces.

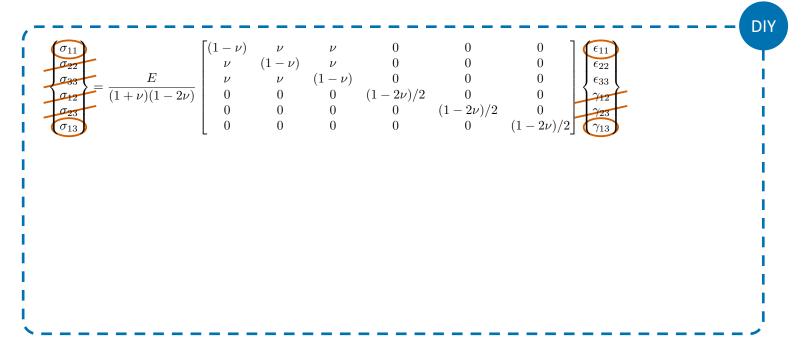
incompatible
with strain field

(Reissner's theory
overcomes this issue)

Timoshenko beams

- Small-strain tensor || Linear elasticity (generalized Hooke's law)

$$\epsilon_{ij} = \begin{bmatrix} \bar{\epsilon}_{11} + z\bar{\kappa}_{11} & 0 & \bar{\gamma}_{13}/2 \\ 0 & 0 & 0 \\ \bar{\gamma}_{13}/2 & 0 & 0 \end{bmatrix} \qquad \sigma_{ij} = \begin{bmatrix} E(\bar{\epsilon}_{11} + z\bar{\kappa}_{11}) & 0 & \frac{E}{2(1+\nu)}\bar{\gamma}_{13} \\ 0 & 0 & 0 \\ \frac{E}{2(1+\nu)}\bar{\gamma}_{13} & 0 & 0 \end{bmatrix}$$



Timoshenko beams

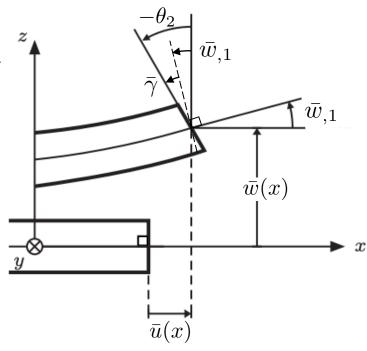
- Linear elasticity (generalized Hooke's law) assuming $\sigma_{22} \ll \sigma_{11}$ and $\sigma_{33} \ll \sigma_{11}$ Since lateral surfaces are traction free, stresses in 22 and 33 directions must be zero on free surfaces.

$$\sigma_{ij} = \begin{bmatrix} E(\bar{\epsilon}_{11} + z\bar{\kappa}_{11}) & 0 & \frac{E}{2(1+\nu)}\bar{\gamma}_{13} \\ 0 & 0 & 0 \\ \frac{E}{2(1+\nu)}\bar{\gamma}_{13} & 0 & 0 \end{bmatrix}$$

Internal resultants

$$N_1 = \int_A \sigma_{11} dA \qquad M_2 = \int_A z \sigma_{11} dA$$
$$V_1 = K_s \int_A \sigma_{13} dA$$

Simplified geometry: $b \ll L$, $h \ll L$



beam geometry

height (h), width (b) cross-sectional area (A) moment of inertia (I) Shear correction factor (K_s)

Timoshenko beams

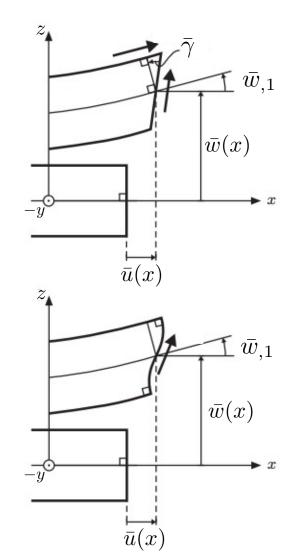
Internal resultants

$$N_{1} = b \int_{-h/2}^{h/2} E(\bar{\epsilon}_{11} + z\bar{\kappa}_{11}) dz = EA\bar{\epsilon}_{11}$$

$$M_{2} = b \int_{-h/2}^{h/2} E(z\bar{\epsilon}_{11} + z^{2}\bar{\kappa}_{11}) dz = EI_{22}\bar{\kappa}_{11}$$

$$V_{1} = K_{s}b \int_{-h/2}^{h/2} \frac{E}{2(1+\nu)} \bar{\gamma}_{13} dz = \frac{K_{s}AE}{2(1+\nu)} \bar{\gamma}_{13}$$

Note: The shear correction factor K_s is required in the formulation due to the absence of shear stress and strain at the top and bottom boundaries of the beam. K_s is computed such that resultant shear force V_1 creates the same strain energy as does the true transverse stresses predicted by the three-dimensional elasticity theory.



Timoshenko beams

- Shear correction factor

$$V_1 = K_{\rm s}b \int_{-h/2}^{h/2} \frac{E}{2(1+\nu)} \bar{\gamma}_{13} dz = \frac{K_{\rm s}AE}{2(1+\nu)} \bar{\gamma}_{13}$$

LXVI. On the Correction for Shear of the Differential Equation for Transverse Vibrations of Prismatic Bars. By Prof. S. P. TIMOSHENKO*.

IN studying the transverse vibrations of prismatic bars, we usually start from the differential equation

$$EI\frac{\partial^4 y}{\partial x^4} + \frac{\rho\Omega}{g}\frac{\partial^2 y}{\partial t^2} = 0, \quad . \quad . \quad . \quad (1)$$

in which EI denotes the flexural rigidity of the bar, Ω the area of the cross-section,

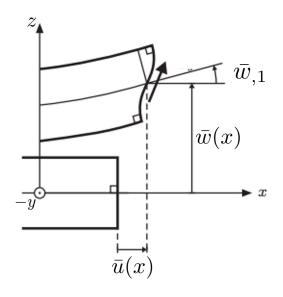
and $\frac{\rho}{\sigma}$ the density of the material.

When the "rotatory inertia" is taken into consideration, the equation takes the form

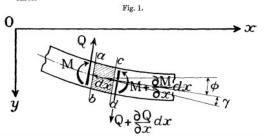
$$EI \frac{\partial^4 y}{\partial x^4} - \frac{I\rho}{g} \frac{\partial^4 y}{\partial x^2 \partial t^2} + \frac{\rho \Omega}{g} \frac{\partial^2 y}{\partial t^2} = 0. \quad . \quad . \quad (2)$$

I now propose to show how the effect of the shear may be taken into account in investigating transverse vibrations, and I shall deduce the general equation of vibration, from which equations (1) and (2) may be obtained as special cases.





S.P. Timoshenko (1921), "On the correction for shear of the differential equation for transverse vibrations of prismatic bars", *Philosophical Magazine and Journal of Science*, 41:245, 744-746.



Dong, S. B., Alpdogan, C., & Taciroglu, E. (2010). "Much ado about shear correction factors in Timoshenko beam theory". International Journal of Solids and Structures, 47(13), 1651-1665.

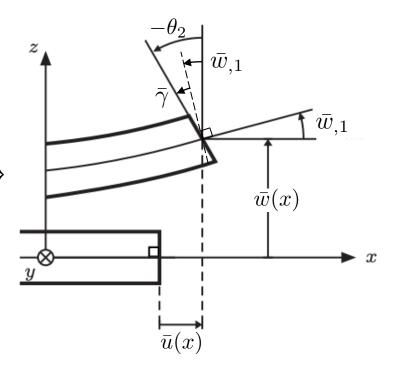
Timoshenko beams

Constitutive relation
 (relationship between internal resultants and kinematic variables)

$$\begin{cases} N_1 \\ M_2 \\ V_1 \end{cases} = \begin{bmatrix} EA & 0 & 0 \\ 0 & EI_{22} & 0 \\ 0 & 0 & \frac{K_{\rm s}EA}{2(1+\nu)} \end{bmatrix} \begin{cases} \bar{\epsilon}_{11} \\ \bar{\kappa}_{11} \\ \bar{\gamma}_{13} \end{cases}$$

Euler-Bernoulli beams

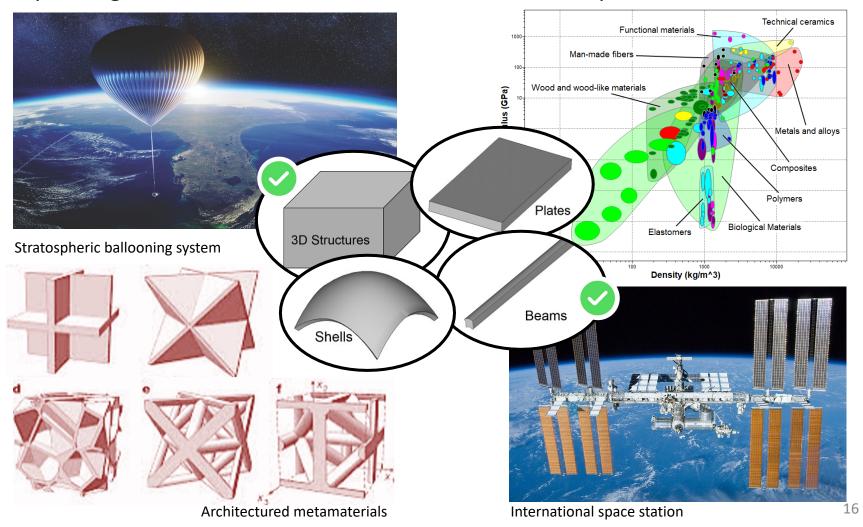
- Constitutive relation



kinematic variables

 $\bar{\epsilon}_{11}(x)$ $\bar{\kappa}_{11}(x)$ $\bar{\gamma}_{13}(x)$ extensional, bending, and shear components of beam strain

Formulation of structural elements (beams, plates, shells) as the analytical upscaling of continuum solids under kinematic assumptions.



Mindlin's plates

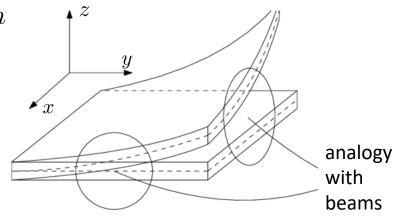
Simplified geometry: the plate is initially flat

- Kinematic assumptions, with $\bar{w} \ll h$

$$u_1(x, y, z) = \bar{u}(x, y) + z\theta_1(x, y)$$

$$u_2(x, y, z) = \bar{v}(x, y) + z\theta_2(x, y)$$

$$u_3(x, y, z) = \bar{w}(x, y)$$



- Small-strain tensor, with $(\bar{w}_{,1})^2 \ll 1$

$$\epsilon_{ij} = \frac{1}{2}(u_{i,j} + u_{j,i})$$

$$\epsilon_{ij} = \begin{bmatrix} \bar{\epsilon}_{11} + z\bar{\kappa}_{11} & \bar{\gamma}_{12}/2 + z\bar{\kappa}_{12} & \bar{\gamma}_{13}/2\\ \bar{\gamma}_{12}/2 + z\bar{\kappa}_{12} & \bar{\epsilon}_{22} + z\bar{\kappa}_{22} & \bar{\gamma}_{23}/2\\ \bar{\gamma}_{13}/2 & \bar{\gamma}_{23}/2 & 0 \end{bmatrix}$$

kinematic variables

$$\bar{\epsilon}_{11}$$
 $\bar{\epsilon}_{22}$ $\bar{\kappa}_{11}$ $\bar{\kappa}_{22}$

extensional strains, bending strains

kinematic variables (cont.)

$$\bar{\gamma}_{12}$$
 $\bar{\kappa}_{12}$ $\bar{\gamma}_{13}$ $\bar{\gamma}_{23}$

contribution of the mid-plane displacement to in-plane shear, contribution of normal rotations to in-plane shear, and shear components

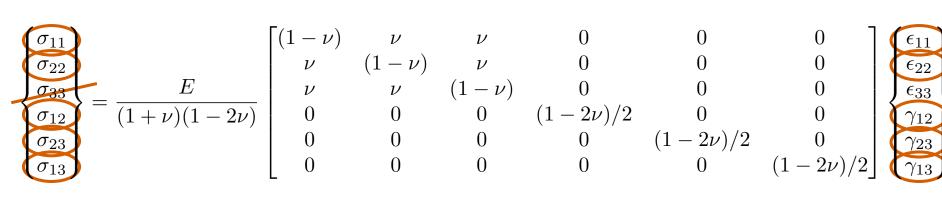
Mindlin's plates

- Small-strain tensor

$$\epsilon_{ij} = \begin{bmatrix} \bar{\epsilon}_{11} + z\bar{\kappa}_{11} & \bar{\gamma}_{12}/2 + z\bar{\kappa}_{12} & \bar{\gamma}_{13}/2 \\ \bar{\gamma}_{12}/2 + z\bar{\kappa}_{12} & \bar{\epsilon}_{22} + z\bar{\kappa}_{22} & \bar{\gamma}_{23}/2 \\ \bar{\gamma}_{13}/2 & \bar{\gamma}_{23}/2 & 0 \end{bmatrix}$$

- Linear elasticity (generalized Hooke's law) assuming $\sigma_{33} \ll \sigma_{22}$ and $\sigma_{33} \ll \sigma_{11}$ Since top/bottom surfaces are traction free, the stress in 33 direction must be zero on free surfaces.

solution of a mixed boundary condition problem



Aside: Kirchhoff plates

- Small-strain tensor, with $\, heta_2 = - ar w_{,1}\,\,,\,\, heta_1 = - ar w_{,2}\,$

$$\epsilon_{ij} = \begin{bmatrix} \bar{\epsilon}_{11} + z\bar{\kappa}_{11} & \bar{\gamma}_{12}/2 + z\bar{\kappa}_{12} & \bar{\gamma}_{13}/2 \\ \bar{\gamma}_{12}/2 + z\bar{\kappa}_{12} & \bar{\epsilon}_{22} + z\bar{\kappa}_{22} & \bar{\gamma}_{23}/2 \\ \bar{\gamma}_{13}/2 & \bar{\gamma}_{23}/2 & 0 \end{bmatrix}$$

- Linear elasticity (generalized Hooke's law) assuming $\sigma_{33} \ll \sigma_{22}$ and $\sigma_{33} \ll \sigma_{11}$ Since top/bottom surfaces are traction free, the stress in 33 direction must be zero on free surfaces.

solution of a mixed boundary condition problem

$$\begin{vmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{12} \\ \sigma_{23} \\ \sigma_{13} \end{vmatrix} = \frac{E}{(1+\nu)(1-2\nu)} \begin{bmatrix} (1-\nu) & \nu & \nu & 0 & 0 & 0 \\ \nu & (1-\nu) & \nu & 0 & 0 & 0 \\ \nu & \nu & (1-\nu) & 0 & 0 & 0 \\ 0 & 0 & 0 & (1-2\nu)/2 & 0 & 0 \\ 0 & 0 & 0 & 0 & (1-2\nu)/2 & 0 \\ 0 & 0 & 0 & 0 & 0 & (1-2\nu)/2 \end{bmatrix} \begin{bmatrix} \epsilon_{11} \\ \epsilon_{22} \\ \epsilon_{33} \\ \gamma_{12} \\ \gamma_{23} \\ \gamma_{13} \\ \gamma_{13} \end{bmatrix}$$

Mindlin's plates

$$\begin{array}{ll} \text{normal} & N_1 = \int_{-h/2}^{h/2} \frac{E}{1-\nu^2} (\bar{\epsilon}_{11} + z\bar{\kappa}_{11} + \nu\bar{\epsilon}_{22} + \nu z\bar{\kappa}_{22}) \mathrm{d}z = \frac{Eh}{1-\nu^2} (\bar{\epsilon}_{11} + \nu\bar{\epsilon}_{22}) \\ & N_2 = \int_{-h/2}^{h/2} \frac{E}{1-\nu^2} (\bar{\epsilon}_{22} + z\bar{\kappa}_{22} + \nu\bar{\epsilon}_{11} + \nu z\bar{\kappa}_{11}) \mathrm{d}z = \frac{Eh}{1-\nu^2} (\bar{\epsilon}_{22} + \nu\bar{\epsilon}_{11}) \\ & \text{in-plane} \\ & \text{shear forces} \end{array} \quad V_{12} = \int_{-h/2}^{h/2} \frac{E}{2(1+\nu)} (\bar{\gamma}_{12} + 2z\bar{\kappa}_{12}) \, \mathrm{d}z = \frac{Eh}{2(1+\nu)} \bar{\gamma}_{12} \\ & \text{bending} \\ & \text{moment} \end{array} \quad M_1 = \int_{-h/2}^{h/2} \frac{E}{1-\nu^2} (\bar{\epsilon}_{11} + z\bar{\kappa}_{11} + \nu\bar{\epsilon}_{22} + \nu z\bar{\kappa}_{22}) z \mathrm{d}z = \frac{Eh^3}{12(1-\nu^2)} (\bar{\kappa}_{11} + \nu\bar{\kappa}_{22}) \\ & \text{flexural state} \\ & M_2 = \int_{-h/2}^{h/2} \frac{E}{1-\nu^2} (\bar{\epsilon}_{22} + z\bar{\kappa}_{22} + \nu\bar{\epsilon}_{11} + \nu z\bar{\kappa}_{11}) z \mathrm{d}z = \frac{Eh^3}{12(1-\nu^2)} (\bar{\kappa}_{22} + \nu\bar{\kappa}_{11}) \\ & \text{twisting} \\ & \text{moment} \end{array} \quad Q_{12} = \int_{-h/2}^{h/2} \frac{E}{2(1+\nu)} (\bar{\gamma}_{12} + 2z\bar{\kappa}_{12}) z \mathrm{d}z = \frac{Eh^3}{24(1+\nu)} \bar{\kappa}_{12} \end{array} \quad N_2 Q_{12}$$

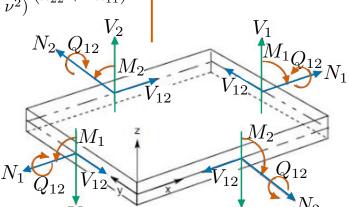
$$Q_{12} = \int_{-h/2}^{h/2} \frac{E}{2(1+\nu)} \left(\bar{\gamma}_{12} + 2z\bar{\kappa}_{12}\right) z dz = \frac{Eh^3}{24(1+\nu)} \bar{\kappa}_1$$

transverse shear forces

$$V_{1} = K_{s} \int_{-h/2}^{h/2} \frac{E}{2(1+\nu)} \bar{\gamma}_{13} dz = \frac{K_{s}Eh}{2(1+\nu)} \bar{\gamma}_{13}$$

$$V_{2} = K_{s} \int_{-h/2}^{h/2} \frac{E}{2(1+\nu)} \bar{\gamma}_{23} dz = \frac{K_{s}Eh}{2(1+\nu)} \bar{\gamma}_{23}$$

$$K_{s} = 5/6$$



Mindlin's plates

Constitutive relation
 (relationship between internal resultants and kinematic variables)

$$\begin{Bmatrix} N_1 \\ N_2 \\ V_{12} \end{Bmatrix} = \frac{Eh}{1 - \nu^2} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{Bmatrix} \bar{\epsilon}_{11} \\ \bar{\epsilon}_{22} \\ \bar{\gamma}_{12}/2 \end{Bmatrix}$$

$$\begin{cases} M_1 \\ M_2 \\ Q_{12} \end{cases} = \frac{Eh^3}{12(1-\nu^2)} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & 1-\nu \end{bmatrix} \begin{bmatrix} \bar{\kappa}_{11} \\ \bar{\kappa}_{22} \\ \bar{\kappa}_{12}/2 \end{cases}$$

$$\begin{Bmatrix} V_1 \\ V_2 \end{Bmatrix} = \frac{K_s E h}{2(1+\nu)} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{Bmatrix} \bar{\gamma}_{13} \\ \bar{\gamma}_{23} \end{Bmatrix}$$

kinematic variables

$$\bar{\epsilon}_{11}$$
 $\bar{\epsilon}_{22}$ $\bar{\kappa}_{11}$ $\bar{\kappa}_{22}$

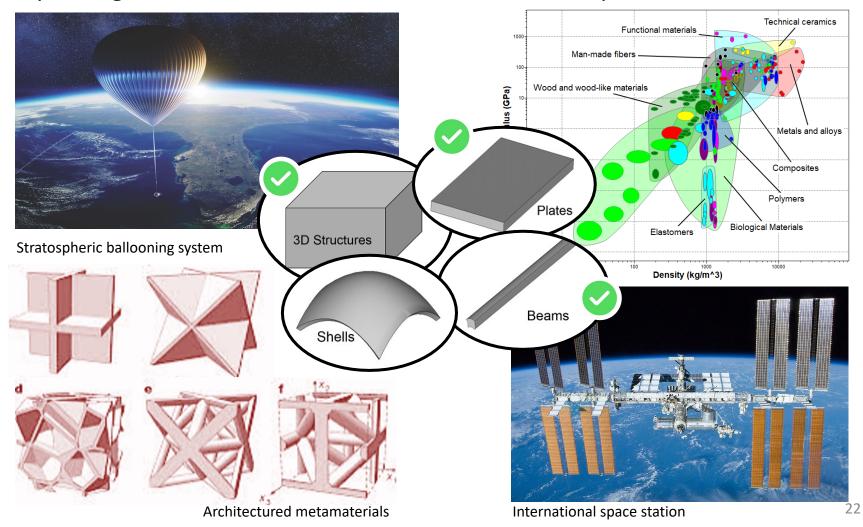
extensional strains, bending strains

kinematic variables (cont.)

$$\bar{\gamma}_{12}$$
 $\bar{\kappa}_{12}$ $\bar{\gamma}_{13}$ $\bar{\gamma}_{23}$

contribution of the mid-plane displacement to in-plane shear, contribution of normal rotations to in-plane shear, and shear components

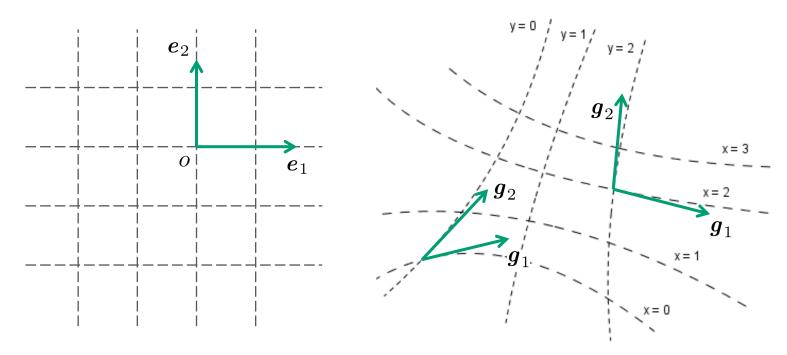
Formulation of structural elements (beams, plates, shells) as the analytical upscaling of continuum solids under kinematic assumptions.



Review Lecture 1: Coordinate system

Curvilinear coordinate system

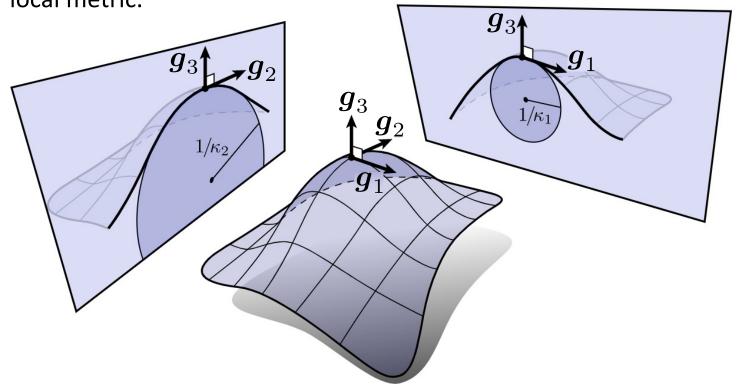
- An <u>origin</u> (relative to which positions are measured)
- A set of coordinate curves
- Basis are defined at each position in space as the tangent vectors $\{g_i\}$ to the coordinate curves. Therefore, basis vectors change from position to position. Basis $\{g_i\}$ are non-orthogonal in general.



Mindlin's shells – Kirchhoff shells

 Like plate formulations but using an orthogonal curvilinear coordinate system (lines of principal curvature on a smooth surface are orthogonal).

Small-strain tensor and kinematic variables now involve local basis and local metric.



Shell buckling



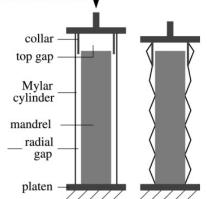






"Yoshimura" pattern

Progressive formation of a surface texture during axial buckling of a thin-walled cylinder fitted onto a mandrel core



applied force

Any questions?