

$$|\psi(t)\rangle = \sum_n c_n e^{-iE_n t} |E_n\rangle$$

BTHP

(1)

Deutsch  
Srednicki

$$\langle \psi(t) | \hat{A} | \psi(t) \rangle = \sum_{n, n'} c_n^* c_n e^{-i(E_n - E_{n'}) t} \langle E_{n'} | \hat{A} | E_n \rangle$$

$$= \sum_n |c_n|^2 \langle E_n | \hat{A} | E_n \rangle + \sum_{n \neq n'} c_n^* c_n e^{-i(E_n - E_{n'}) t} \langle E_{n'} | \hat{A} | E_n \rangle$$

$\langle E_{n'} | \hat{A} | E_n \rangle$  small  $\rightarrow$  fluctuations.

$c_n$ : off-diag. elements odd exp.

$$\langle \psi(t) | \hat{A} | \psi(t) \rangle \Big|_{t \gg t_{\text{mix}}} \approx \sum_n |c_n|^2 \langle E_n | \hat{A} | E_n \rangle$$

$$\langle E_n | \hat{A} | E_n \rangle = A(E_n)$$

$$c_n \sim 1/\sqrt{\Omega}$$

# of states

$$\sum_{n, n'} c_n^* c_n \sim 1$$

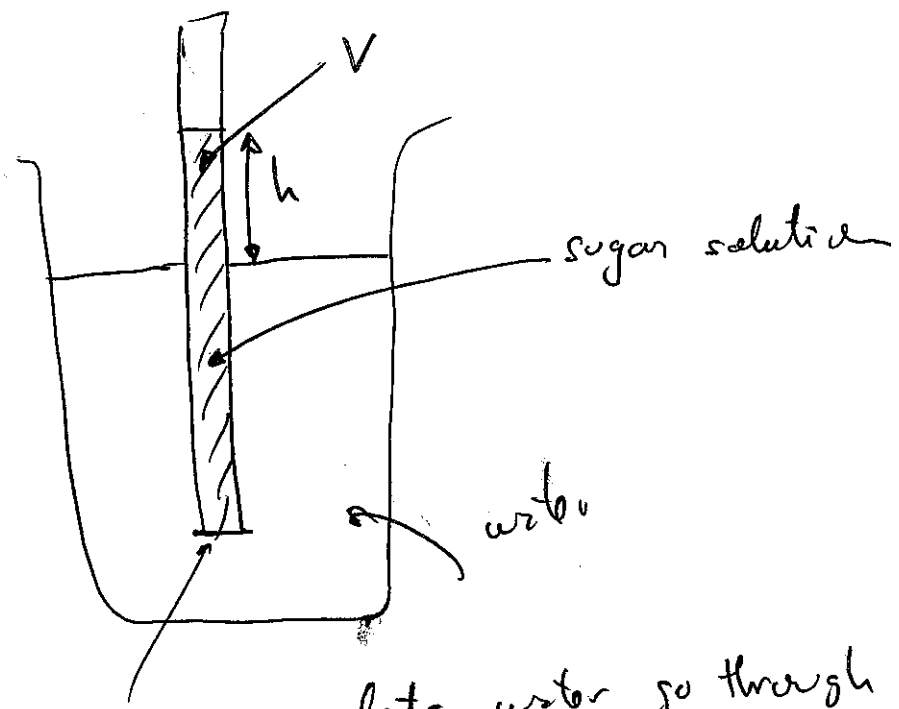
$$\langle E_n | \hat{A} | E_n \rangle$$

random phases  $\sim$  random walk

$$\sqrt{\Omega} \sim \Omega$$

$$\sim 1/\sqrt{\Omega} \text{ avg.}$$

# Osmotic pressure



membrane, lets water go through but not sugar.

$$p = \rho g h = \frac{n_1 R T}{V} = \frac{N_1 k_B T}{V}$$

$n_1$ : moles of solute

thermodynamics of weak solutions.

$$c = \frac{n_{\text{solute}}}{N_{\text{solvent}}} \quad c \ll 1.$$

$\Phi_0(P, T, N)$  Gibbs potential

①

$$\Phi = E - TS + pV = \mu N$$

$$E = TS - pV + \mu N.$$

$$d\Phi = -SdT + Vdp + \mu dN$$

$$\Phi_0 = N\mu_0(P, T)$$

add one molecule of solute.

$$\Delta\Phi = \alpha(P, T, N)$$

non-interacting  $\Rightarrow$  add  $n \rightarrow \Delta\Phi = n\alpha(P, T, N)$

However we have to divide by  $n!$  (identical particles).

$$-\beta\Delta \rightarrow -\ln n! = -n \ln n + n$$

$$\Delta \rightarrow \Delta + k_B T (n \ln n - n)$$

$$\Phi \rightarrow \Phi + n\alpha + k_B T (n \ln n - n)$$

$$\Phi = N\mu_0(P, T) + n\alpha(P, T, N) + k_B T n \ln(n/e)$$

$$= N\mu_0(P, T) + nk_B T \ln\left(\frac{n}{e} e^{\alpha/k_B T}\right)$$

homogeneous in  $n, N$

$$e^{\alpha/T} = \frac{f(P,T)}{N}$$

$$\Rightarrow \Phi = N\mu_0(P,T) + n k_B T \ln\left(\frac{n}{eN} f(P,T)\right)$$

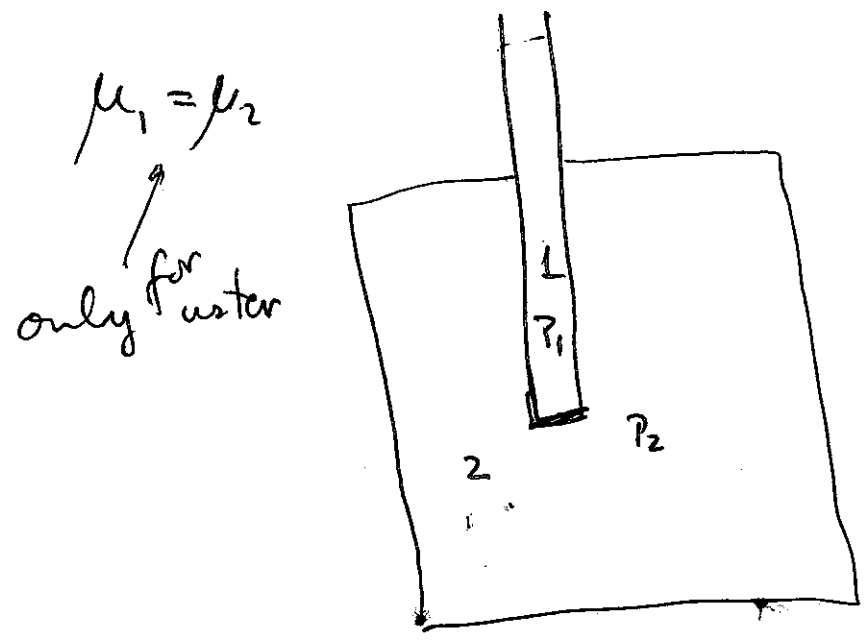
$$= N\mu_0(P,T) + n k_B T \ln\left(\frac{n}{eN}\right) + \underbrace{n k_B T \ln f(P,T)}_{\psi(P,T)}$$

$$\Phi = N\mu_0(P,T) + n \psi(P,T) + n k_B T \ln\left(\frac{n}{eN}\right) + \mathcal{O}\left(\frac{n^2}{N}\right)$$

$$\mu = \frac{\partial \Phi}{\partial N} = \mu_0(P,T) - \frac{n}{N} k_B T ; \mu \text{ of solvent.}$$

$$\begin{aligned} \mu' = \frac{\partial \Phi}{\partial n} &= \psi(P,T) + k_B T \ln\left(\frac{n}{eN}\right) - k_B T \\ &= k_B T \ln\left(\frac{n}{N}\right) + \psi(P,T) \end{aligned}$$

Osmotic pressure :



$$\mu_0(P_1, T) - c \frac{n}{N} k_B T = \mu_0(P_2, T)$$

$$P_1 \approx P_2 \Rightarrow P_1 = P_2 + \delta P$$

$$\frac{\partial \mu_0}{\partial P} \cdot \delta P = c k_B T$$

$$\delta P = \frac{c k_B T}{\partial \mu_0 / \partial P}$$

$$= \frac{n}{N} \frac{k_B T}{V/N} = \frac{n}{V} k_B T \rightarrow$$

$$N \frac{\partial \mu}{\partial P} = V$$

$$\frac{\partial \mu}{\partial P} = V/N$$

$$\delta P_{\text{osmotic}} = \frac{n}{V} k_B T \quad \checkmark$$

From Einstein paper.


Should be same for a suspension of small particles?

$$-\beta A = \ln \int d^{3N} x d^{3N} p d^{3N} \phi e^{-\frac{1}{k_B T} E}$$

E-TS

$$S = \frac{E}{k_B T} + k_B \ln \int d^3 \dots e^{-\beta E}$$

$V^*$ : volume of solution + solvent.

 position of particle 1.

$$Z = \int d^3 x_1 \dots d^3 x_n \underbrace{\int d^3 \dots e^{-\beta E}}_{\text{indep. of } V^*}$$

indep. of  $V^*$

because no water does not see the wall

also indep. of  $\vec{x}_1, \dots, \vec{x}_n$

because of translation symmetry

$$Z = (V^*)^n \cdot Z_0$$

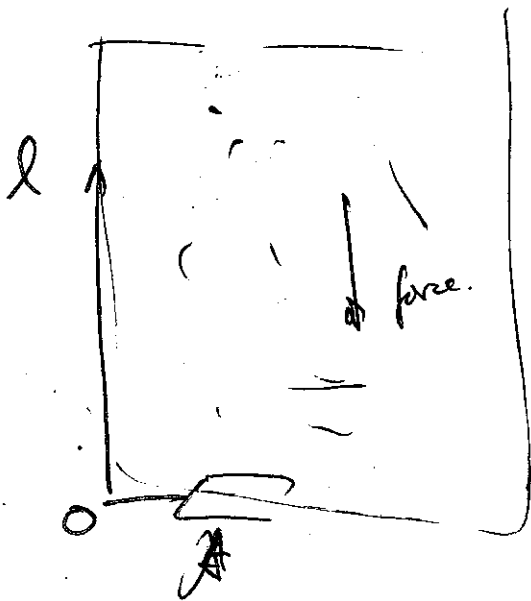
$$-\beta A = n \ln V'' + \ln Z_0$$

$$A = -nk_B T \ln V'' - k_B T \ln Z_0$$

$$dA = -SdT - pdV + \mu dn$$

$$\Phi = -\frac{\partial A}{\partial V''} = \frac{nk_B T}{V''} \quad \checkmark$$

necessary to move wall.



Suspended particles.  $v = n/V$

$$\delta A = \delta E - T\delta S = 0 \quad \text{equilibrium.}$$

$V(x)$  : density of particles

(8)

$$E = A \int_0^l dx v(x) \underbrace{mgx}_{\text{gravity}}$$

$$S = ?$$

From dilute solution

$$\Phi = N\mu_0(p,T) + n \phi(p,T) + nk_B T \ln(n/e\alpha)$$

$$S = - \frac{\partial \Phi}{\partial T} = -N \frac{\partial \mu_0}{\partial T} - n \frac{\partial \phi}{\partial T} + nk_B \ln(n/e\alpha)$$

For a small portion of solution  $\int dx$

$$\Delta S = N \Delta S_0 = \underbrace{\rho \cdot A \cdot dx}_{\text{solvent}} S_0 + A \cdot dx v(x) \overset{\uparrow}{S} \text{ of solute}$$

$$= A \cdot dx k_B v(x) \ln \left( \frac{v(x)}{e\rho} \right)$$

↳  $k_B$  of solvent.

$$S = \underbrace{S_0}_{\text{solvent}} + \underbrace{n S}_{\substack{\uparrow \\ \text{of solute} \\ \text{indep. of } v(x)}} = A k_B \int v(x) \ln \left( \frac{v(x)}{e\rho} \right) dx$$



Free energy:

(9)

$$A = A \left\{ \int_0^{\ell} dx v(x) mg dx + k_B T \int_0^{\ell} v(x) \ln \left( \frac{v(x)}{e\rho} \right) dx + \text{const} \right\}$$

Lagrange multiplier

$$A = A \left\{ \int_0^{\ell} dx v(x) mg dx + k_B T \int_0^{\ell} v(x) \ln \left( \frac{v(x)}{e\rho} \right) dx + \lambda \left( \int_0^{\ell} v(x) dx - \frac{\eta}{A} \right) \right\}$$

$$\frac{\delta A}{\delta v(x)} = 0$$

$$mgx + k_B T \ln \frac{v(x)}{e\rho} + k_B T + \lambda = 0$$

$$mgx + k_B T \ln v(x) + (k_B T \ln(e\rho) + k_B T) + \lambda = 0$$

$$\ln v(x) = -\frac{mgx}{k_B T} + \ln(e\rho) + \frac{\lambda}{k_B T} \Rightarrow v(x) = C e^{-\frac{mgx}{k_B T}}$$

$$A \int_0^l v(x) dx = CA \int_0^l e^{-\frac{mgx}{k_B T}} dx = \quad (10)$$

$$= CA \left. \frac{e^{-\frac{mgx}{k_B T}}}{-\frac{mg}{k_B T}} \right|_0^l = -\frac{CA k_B T}{mg} \left( e^{-\frac{mgl}{k_B T}} - 1 \right)$$

$\approx 0$

$$= \frac{CA k_B T}{mg} = n \Rightarrow C = \frac{mg n}{A k_B T}$$

$$v(x) = \frac{mg}{A k_B T} n e^{-\frac{mgx}{k_B T}}$$

$$P_{\text{osmotic}} = v k_B T = \frac{mg}{A} n e^{-\frac{mgx}{k_B T}}$$

$$\frac{\partial P}{\partial x} = -\frac{mg}{k_B T} \frac{mg}{A} n e^{-\frac{mgx}{k_B T}} = -mg v(x)$$

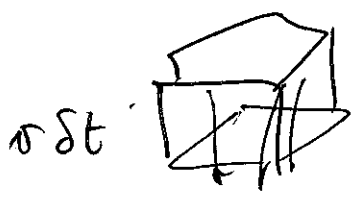
↑  
pressure gradient

↗  
force.

o spherical particles moving in fluid

terminal velocity  $v = \frac{mg}{6\pi\eta R}$

↑ radius  
friction



$v \times \sigma \delta t \times A$

flux  $\frac{vmg}{6\pi\eta R}$

Diffusion  $-D \frac{\partial v}{\partial x}$  flux created by gradient of concentration.  $v$

↑  
coeff. of diffusion

$c = \frac{n}{N} = \frac{v}{p}$

↑ gradient of  $v$  ↓ flux  $-D \frac{\partial v}{\partial x} = \frac{vmg}{6\pi\eta R}$

$\frac{\partial v}{\partial x} = -\frac{mg}{k_B T} v$

$D = \frac{k_B T}{6\pi\eta R}$

$D \frac{mg}{k_B T} v = \frac{vmg}{6\pi\eta R}$

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In Einstein's paper gas constant  $k_B = R/N_A$  Avogadro's #

$$D = \frac{k_B T}{6\pi\eta R}$$

$\eta$   
diffusion.

$R$   
friction

↔  
relativa.