ECE 595, Section 10 Numerical Simulations Lecture 14: Beam Propagation Method

Prof. Peter Bermel February 11, 2013

Outline

- Recap from Friday
- Derivation of Beam Propagation Method
- Nonlinear Schrodinger equation
- Comparison of BPM Strategies

— FFT

- Uniform spatial grid
- Finite element
- Perfectly Matched Layers

Recap from Friday

- Rationale for FFTW
- Planning DFTs
- Executing DFTs
 - Basic interface
 - Advanced interface
- BPM example

Recap: Beam Propagation

• Starting from the Helmholtz equation: m(x) > 2

$$-\nabla^2 \psi = \left(\frac{n\omega}{c}\right)^2 \psi$$

- One can assume a solution of the form: $\psi = \phi e^{-j\beta z}$
- Where ϕ is slowly varying, which gives rise to: $-\nabla^2 \phi + 2j\beta \hat{z} \cdot \nabla \phi = k_{\perp}^2 \phi$

Beam Propagation

 To simplify problem, drop second derivatives in z – now we can write as:

$$\frac{\partial \phi}{\partial z} = \frac{j}{2\beta} \nabla_{\perp}^{2} \phi + \frac{jk_{\perp}^{2}}{2\beta} \phi$$

• Can simplify by defining two operators:

$$U = \frac{j}{2\beta} \nabla_{\perp}^{2}$$
$$W = \frac{jk_{\perp}^{2}}{2\beta}$$
$$\frac{\partial \phi}{\partial z} = (U + W)\phi$$

ECE 595, Prof. Bermel

Nonlinear Schrodinger Equation

 Can derive expressions suitable for understanding fibers with dispersion and Kerr nonlinearity:

$$U = -\frac{j\beta_2}{2}\frac{\partial^2}{\partial t^2}$$
$$W = -\alpha + \frac{j\kappa}{2}|\phi|^2$$
$$\frac{\partial\phi}{\partial z} = (U+W)\phi$$

Nonlinear Schrodinger Equation

- In the presence of nonlinearity, don't actually know the value of W(z+h)
- Can obtain the result iteratively

- Use W(z) to evaluate W(z+h)

- Work backwards to refine guess for $\phi(z+h)$
- After a few iterations, generally reach a selfconsistent solution

Beam Propagation

• For a small z-step of size *h*, we can formally write a solution:

$$\phi(z+h) = e^{h(U+W)}\phi(z)$$

• If we know that U and W operators commute, we can rewrite as:

$$\phi(z+h) = e^{hU}e^{hW}\phi(z)$$

$$\phi(z+h) = e^{hU/2}e^{hW}e^{hU/2}\phi(z)$$

Beam Propagation

- Split-step method
 - Propagate half a step with the Laplacian
 - Propagate linear phase shift over the full distance
 - Propagate half a step with the Laplacian



BPM Strategies

- Most important decision is handling inhomogeneity well
- Possible strategies:
 - FFT
 - Uniform spatial grid
 - Finite-element method

FFT BPM

• Well-suited for diffraction step, where we can rephrase the operator as:

$$U = -\frac{j}{2\beta}(k+G)_{\perp}^{2}$$

• Can transform before and then back afterwards, via FFT

Uniform Spatial Grid BPM

- Reformulate Laplacian in 2D with: $\nabla^2 \phi \approx \frac{\phi_{i-N} + \phi_{i-1} - 4\phi_i + \phi_{i+1} + \phi_{i+N}}{h^2}$
- Where *h* is the grid spacing

Finite Element BPM

- Finite element method consists of dividing a spatial domain in 1D, 2D or 3D into a mesh
- Mesh generally has D+1 vertices
- Solution can take various forms, but usually a tent function within each D+1-gon



ECE 595, Prof. Bermel

Finite Element BPM

• In general, can formulate FE problems as:

Lu = b

- L is the stiffness matrix, representing overlap between basis functions
- b is the integral of given PDE with respect to basis
- *u* is unknown
- Value of FEA comes from:
 - *Spatial flexibility*: can define each element to vary in size quite substantially
 - Speed: properly chosen basis functions have compact support, leading to a sparse matrix

Perfectly Matched Layers

- In order to prevent lateral reflections (e.g., from PEC boundaries), can introduce perfectly matched layers (PML)
- Several formulations (including split-field and uniaxial), but here we'll follow stretched coordinate PML
- Effected by the transformation:

$$\nabla \rightarrow A \cdot \nabla$$

$$A = \begin{pmatrix} 1 - j\beta & 0 & 0 \\ 0 & 1 - j\beta & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\beta = -\frac{3\lambda\rho^2}{4\pi n d^3} \ln R$$

ECE 595, Prof. Bermel

Next Class

- Is on Wednesday, Feb. 13
- Will continue with beam propagation method
- Recommended reading: Obayya, Sections 2.7-2.8